

The horn, the hadron mass spectrum and the QCD phase diagram

(hadron production at chemical freeze-out in central nucleus-nucleus collisions)

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- The thermal model and the thermal fits in (central) AA collisions
- Thermal fits and the QCD phase diagram
- Thermal model and heavy-flavored hadrons (in AA and elementary collisions)

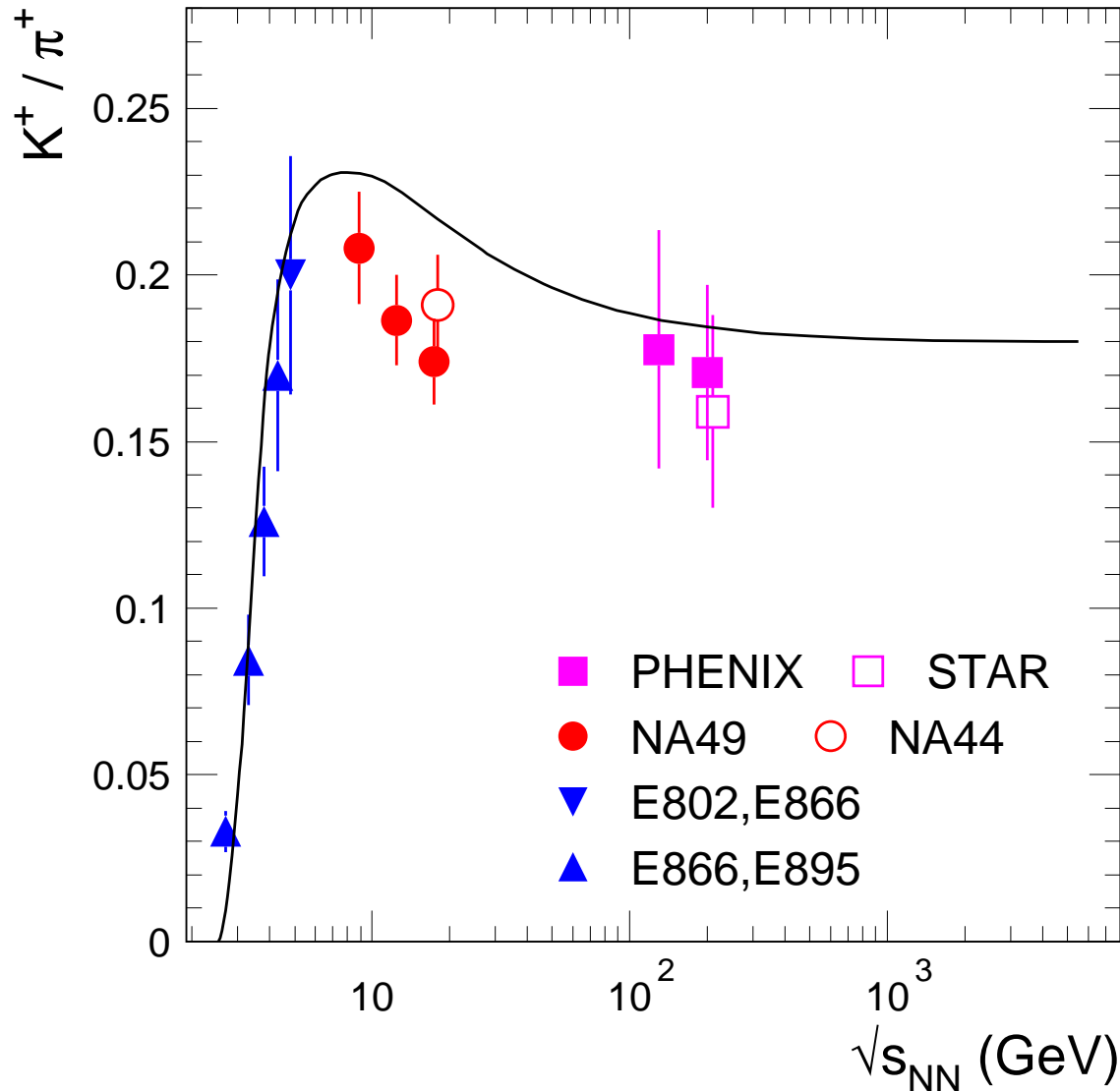
AA, P.Braun-Munzinger, J.Stachel, Phys. Lett. B 673 (2009) 142

AA, P.Braun-Munzinger, K.Redlich, J.Stachel, Phys. Lett. B 652 (2007) 259; B 678 (2009) 350; arXiv:1002.4441

The horn

... as of 2005 ↘

...not well reproduced by the thermal model (line)

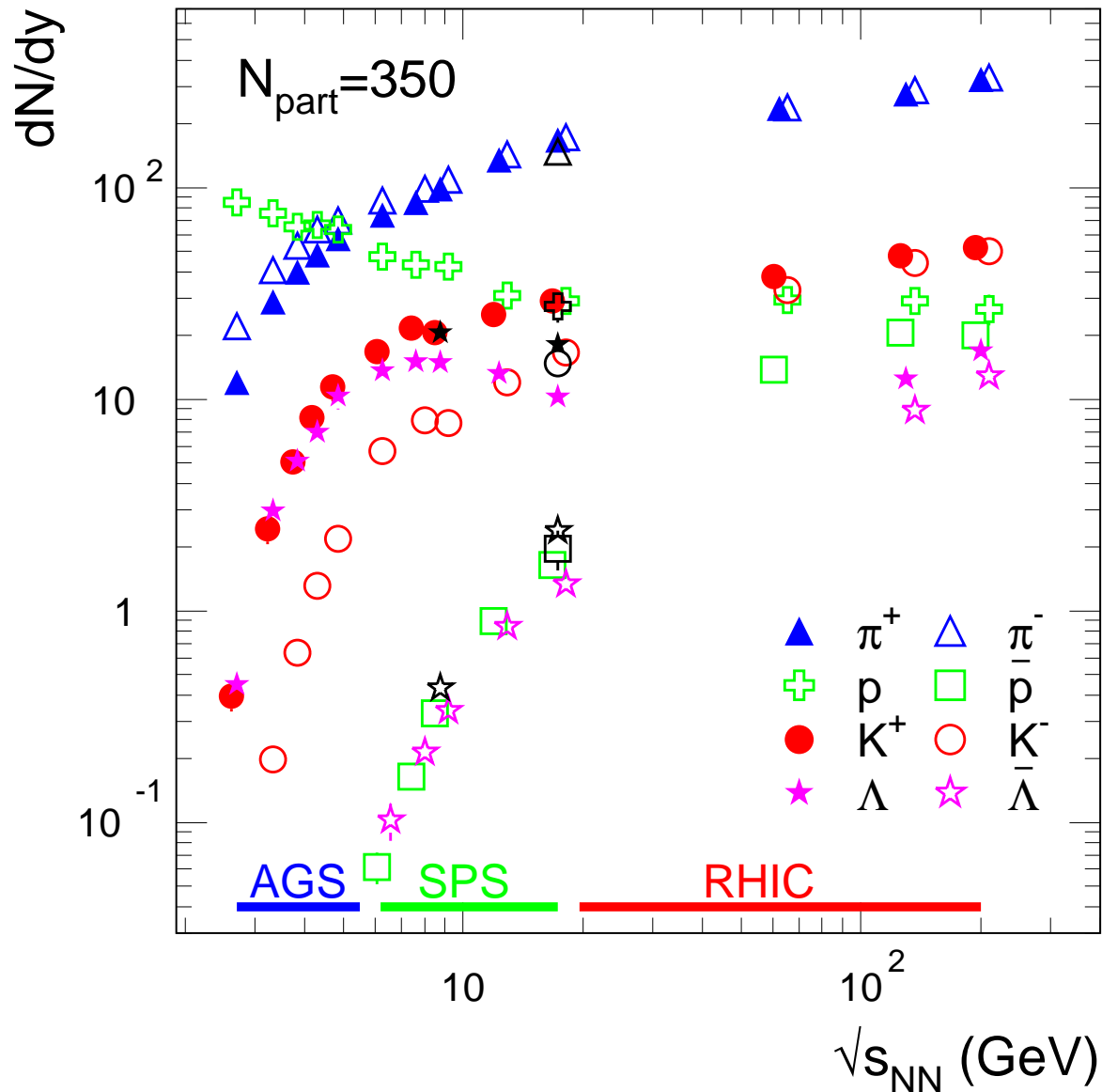


(the 2008 version) taken as experimental evidence for the onset of deconfinement and quark-gluon plasma formation

NA49 collab., Phys. Rev. C **77** (2008)

...as predicted by Gaździcki and Gorenstein, Acta Phys. Polon. B **30** (1999) 2705

Hadron yields 2008 (central collisions)



- lots of particles, mostly newly created ($m = E/c^2$)
- a great variety of species:
 - π^\pm ($u\bar{d}$, $d\bar{u}$), $m=140$ MeV
 - K^\pm ($u\bar{s}$, $\bar{u}s$), $m=494$ MeV
 - p (uud), $m=938$ MeV
 - Λ (uds), $m=1116$ MeV
 - also: $\Xi(dss)$, $\Omega(sss)$...
- mass hierarchy in production (u, d quarks: remnants from the incoming nuclei)

The hadron mass spectrum as of 2008

Particle Data Group, Phys. Lett. B 667 (2008) 1

Additions (compared to 2005):

Many new resonances up to 3 GeV

+(86)4 (non)strange mesons

+(36)30 (non)strange baryons

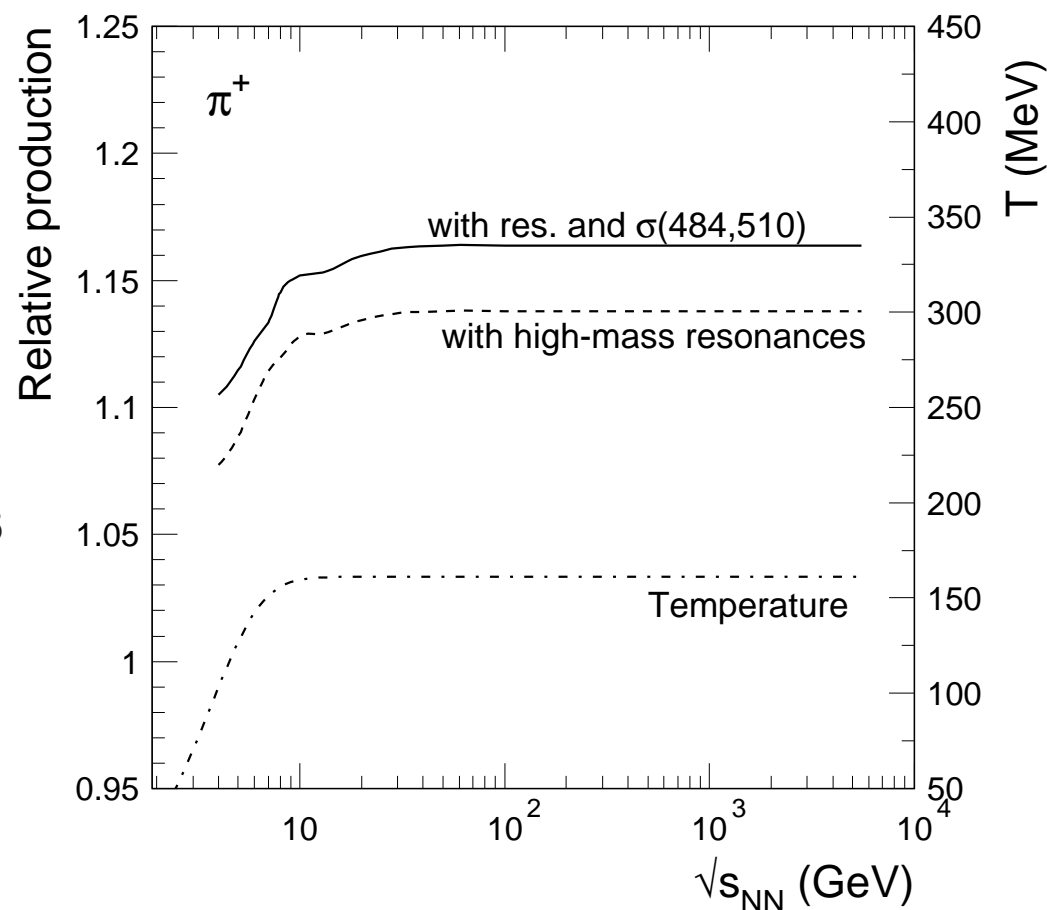
σ meson ($f_0(600)$):

$m_\sigma = 484 \pm 17$ MeV, $\Gamma_\sigma = 510 \pm 20$ MeV

García-Martín, Peláez, Ynduráin, Phys. Rev. D 76
(2007) 074034

(in total 485 hadron species, incl. composites)

relative increase of calc. dens.



The thermal model

grand canonical partition function for specie i ($\hbar = c = 1$):

$$\ln Z_i = \frac{V g_i}{2\pi^2} \int_0^\infty \pm p^2 dp \ln[1 \pm \exp(-(E_i - \mu_i)/T)]$$

$g_i = (2J_i + 1)$ spin degeneracy factor; T temperature;

$E_i = \sqrt{p^2 + m_i^2}$ total energy; (+) for fermions (-) for bosons

$$\mu_i = \mu_b B_i + \mu_{I_3} I_{3i} + \mu_S S_i + \mu_C C_i$$

μ ensure conservation (on average) of quantum numbers:

i) baryon number: $V \sum_i n_i B_i = N_B$

ii) isospin: $V \sum_i n_i I_{3i} = I_3^{tot}$

iii) strangeness: $V \sum_i n_i S_i = 0$

iv) charm: $V \sum_i n_i C_i = 0$.

Short-range repulsive core modelled via excluded volume correction (Rischke)

Widths of resonances taken into account

The thermal fits

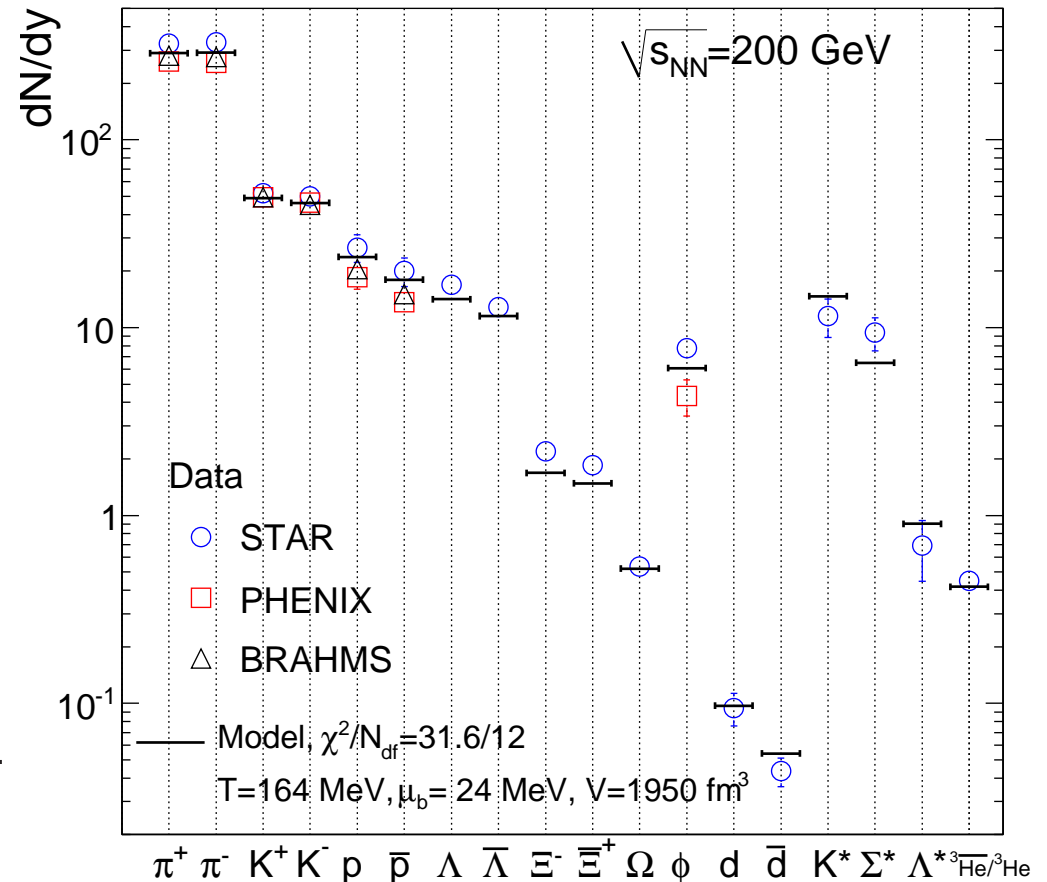
$$n_i = N_i/V = -\frac{T}{V} \frac{\partial \ln Z_i}{\partial \mu} = \frac{g_i}{2\pi^2} \int_0^\infty \frac{p^2 dp}{\exp[(E_i - \mu_i)/T] \pm 1}$$

Minimize: $\chi^2 = \sum_i \frac{(N_i^{exp} - N_i^{therm})^2}{\sigma_i^2}$

N_i : hadron yield ($\Rightarrow T, \mu_b, V$) or yield ratio (no V)

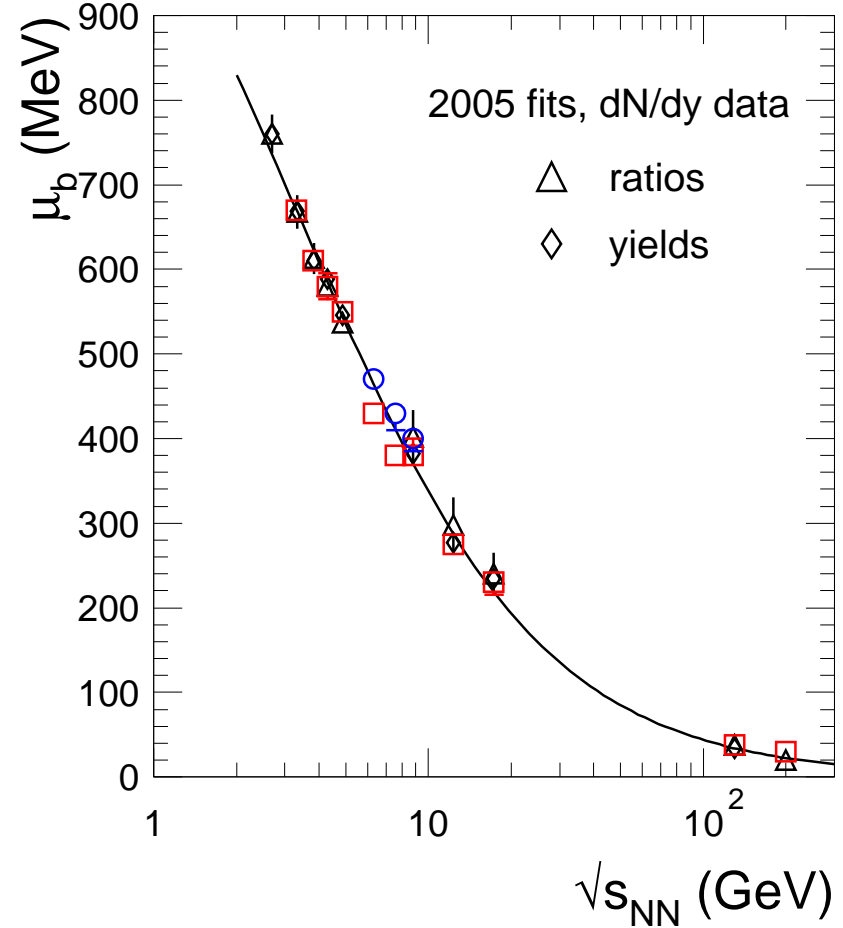
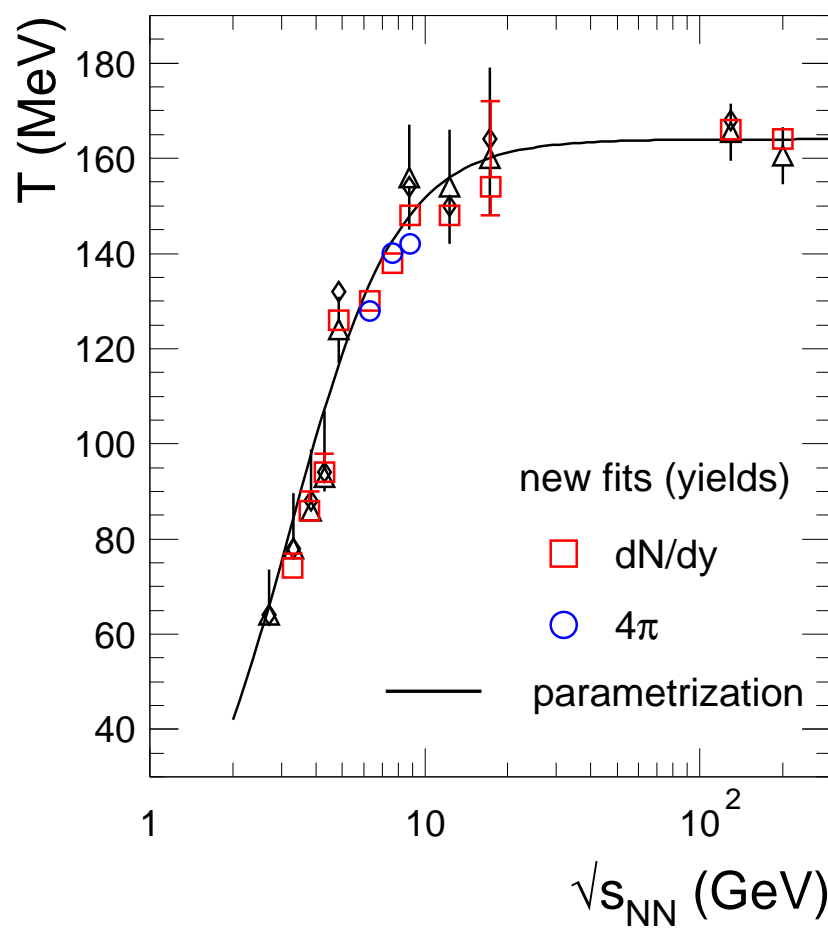
Data: 4π or dN/dy (at $y=0$)

only STAR data: $T=162$ MeV, $\mu_b=24$ MeV, $V=2400$ fm³, $\chi^2/N_{df}=10.9/12$



The hadron abundances are in agreement with a thermally equilibrated system

Energy dependence of T , μ_b



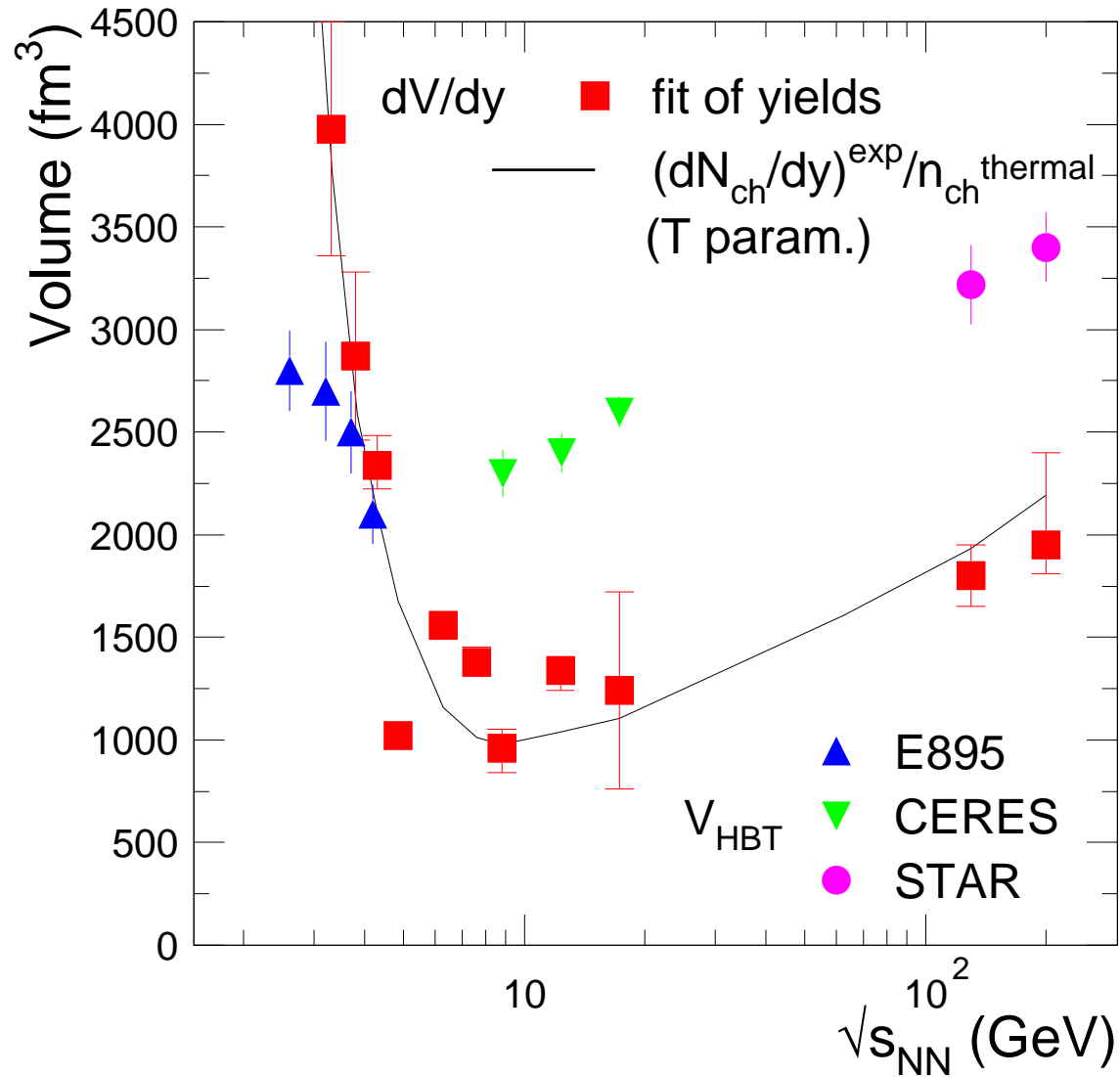
thermal fits exhibit a limiting temperature:

$$T = T_{lim} \frac{1}{1 + \exp(2.60 - \ln(\sqrt{s_{NN}}(\text{GeV}))/0.45)},$$

$$T_{lim} = 164 \pm 4 \text{ MeV}$$

$$\mu_b[\text{MeV}] = \frac{1303}{1 + 0.286 \sqrt{s_{NN}}(\text{GeV})}$$

Energy dependence of the freeze-out volume



dV/dy : volume for one unit rapidity (at midrapidity)

minimum at $T \rightarrow T_{\text{lim}}$

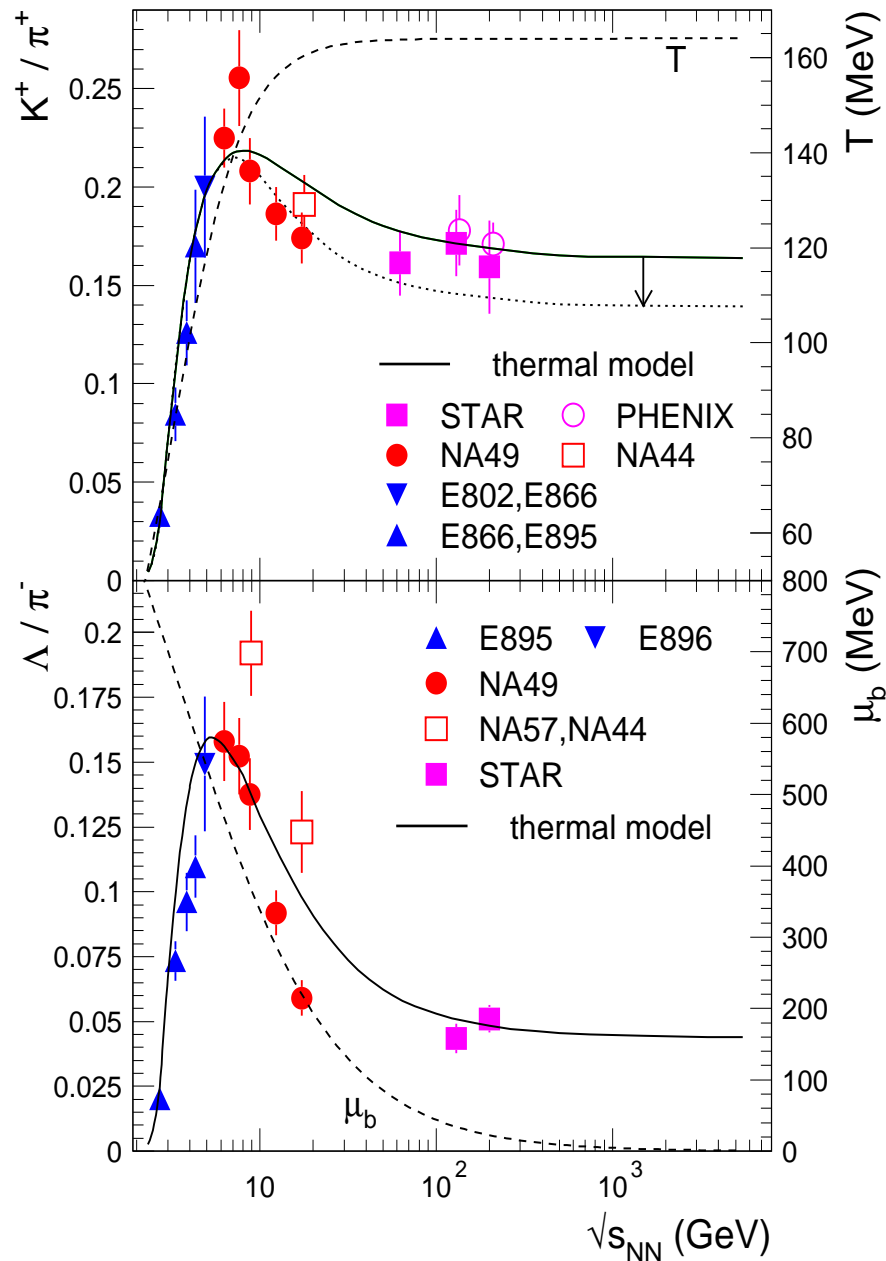
V_{HBT} :

CERES, PRL, 90 (2003) 022301

($\lambda_f \simeq 1 \text{ fm}$)

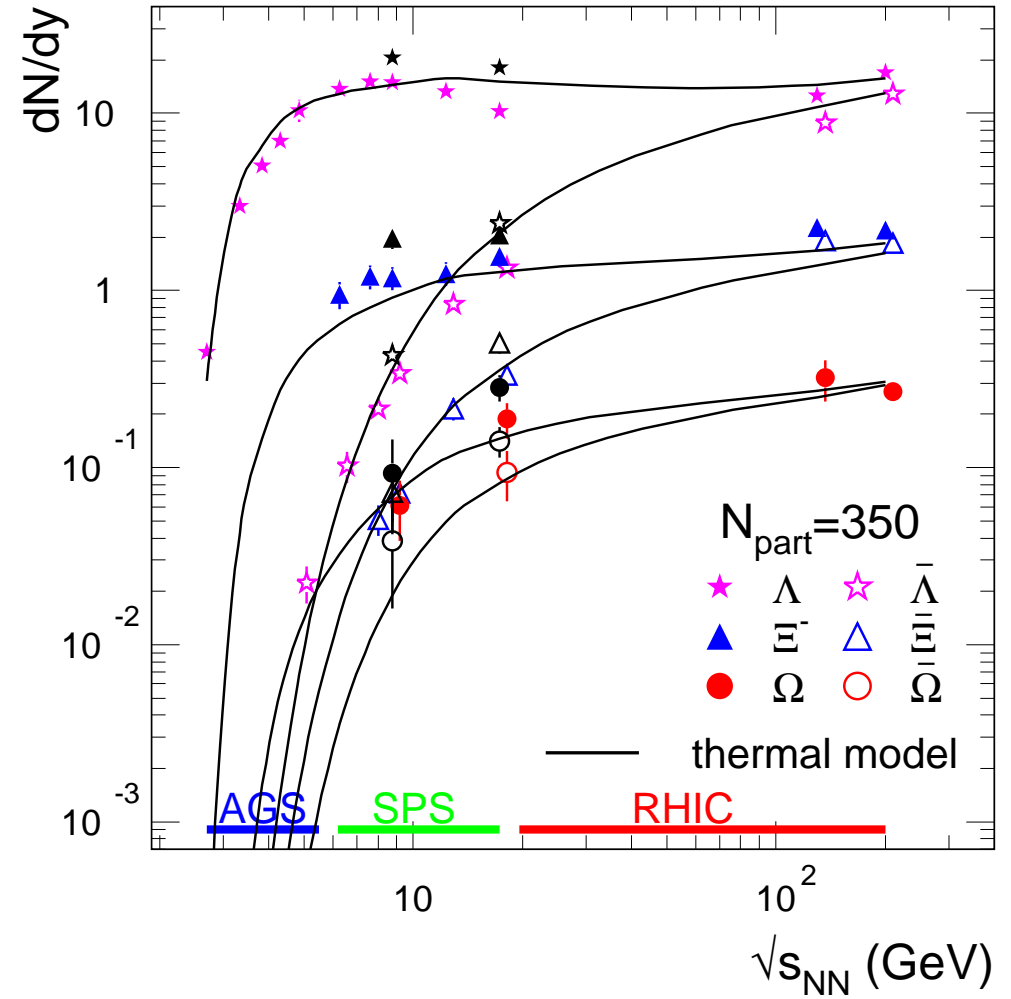
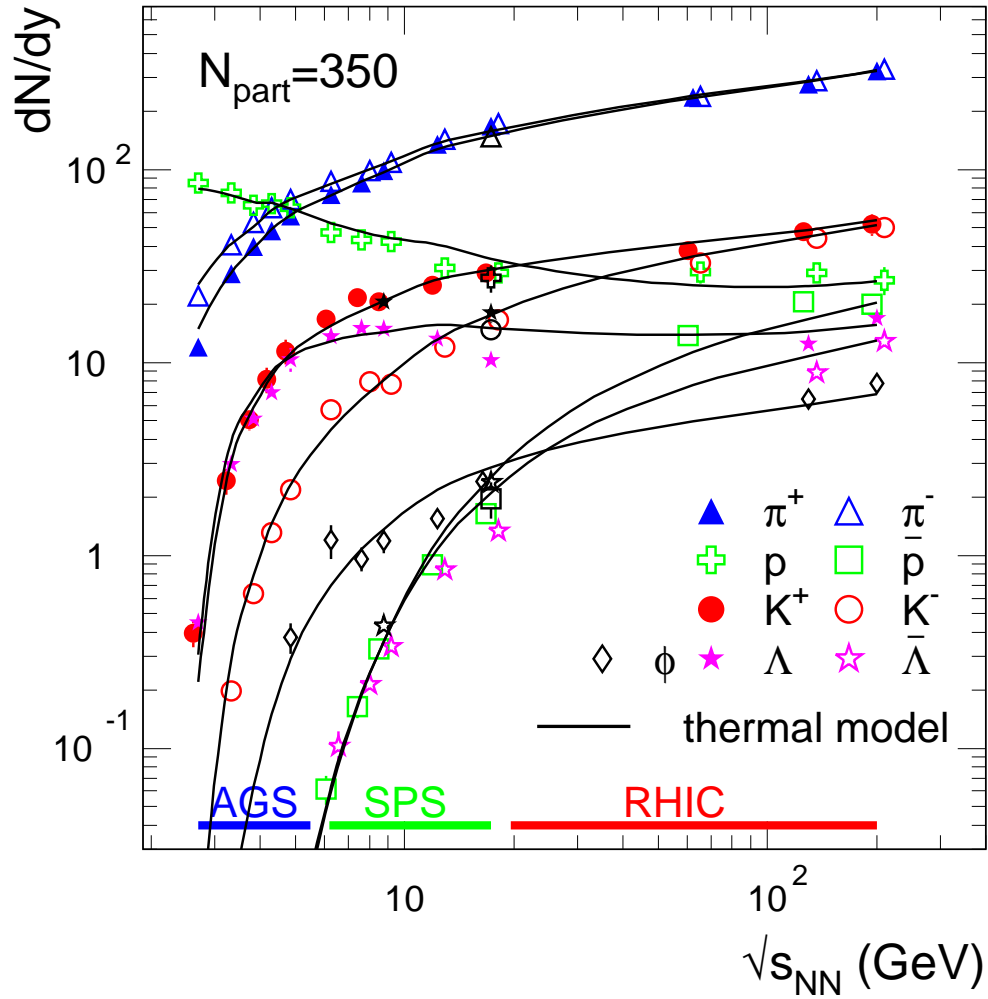
not fully understood dependence

The horn as of 2009

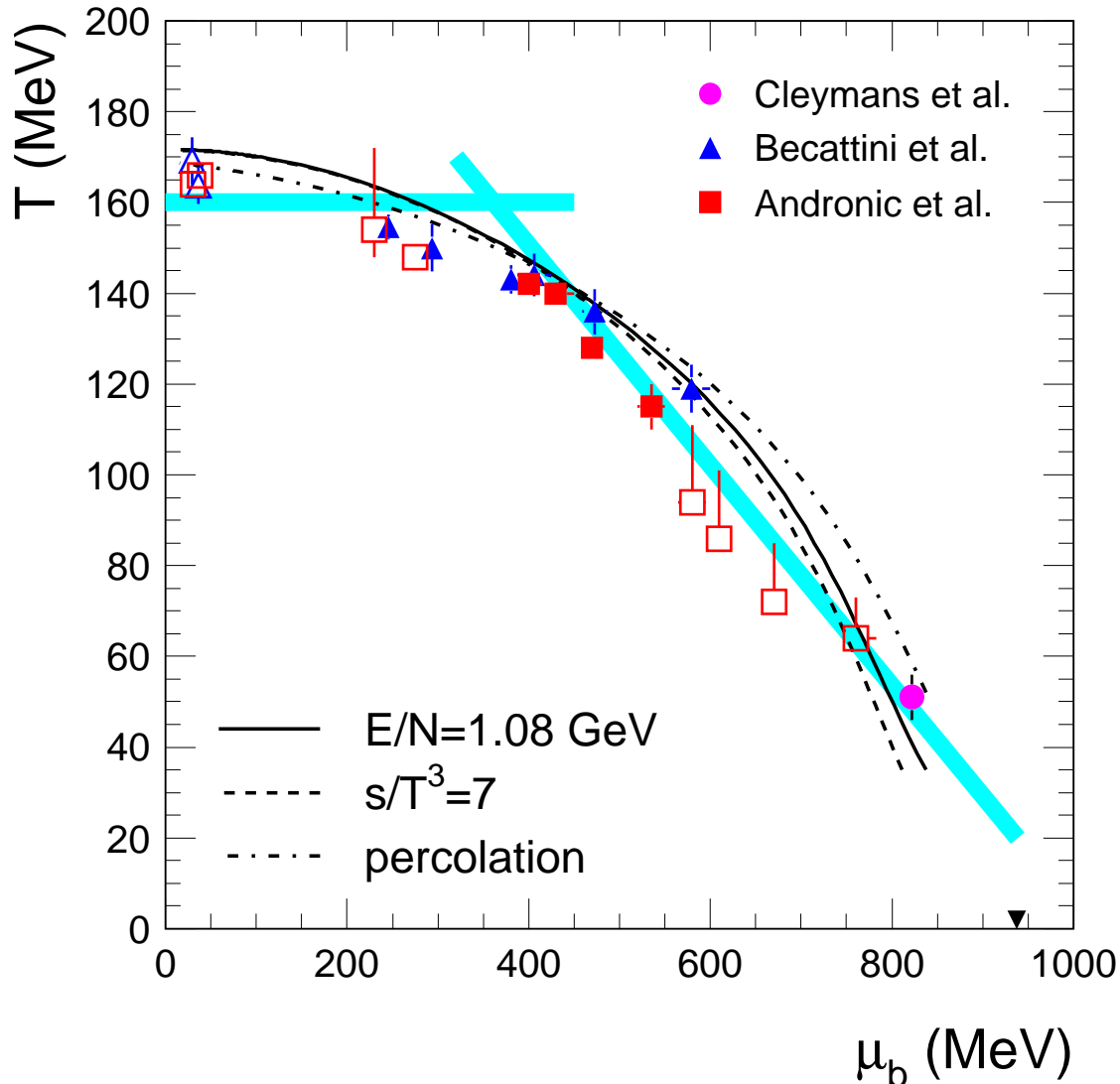


much better explained by the model
 ...as due to detailed features of the hadron mass spectrum
 ...which leads to a limiting temperature (“Hagedorn”, $T < T_H$)
 ...and contains the QCD phase transition
 the horn’s sensitivity to the phase boundary is determined (via strangeness neutrality condition) by the Λ abundance (determined by both T and μ_b)

Yields at mid-rapidity



The phase diagram of QCD



is chemical freeze-out a determination of the phase boundary?
if yes, how is thermalization achieved?

- for SPS energies and higher:
driven by the deconfinement transition

PBM, Stachel, Wetterich, PLB 596 (2004) 61

- for lower energies (SIS100):
is the quarkyonic phase transition the “thermalizer”?

McLerran, Pisarski, NPA 796 (2007) 83; see also: AA et al., NPA 837 (2010) 65

Statistical model for heavy quark hadrons

P.Braun-Munzinger, J.Stachel, PLB 490 (2000) 196

Assumptions

- all charm quarks are produced in primary hard collisions ($t_{c\bar{c}} \sim 1/2m_c \simeq 0.1 \text{ fm}/c$)
- survive and thermalize **in QGP** (thermal, but not chemical equilibrium)
- charmed hadrons are formed at chemical freeze-out together with all hadrons
statistical laws, quantum numbers conservation
statistical hadronization (\neq coalescence)
is freeze-out at/the(?) phase boundary? can we delineate it more with charm?
- no J/ψ survival in QGP (full screening) ...can J/ψ survive above T_c ? (LQCD)

Asakawa, Hatsuda, PRL 92 (2004) 012001; Mocsy, Petreczky, PRL 99 (2007) 211602

Timescales for charm(onium) production

Karsch & Petronzio, PLB 193 (1987) 105, Blaizot & Ollitrault, PRD 39 (1989) 232

- QGP formation time, t_{QGP}
 - SPS (FAIR): $t_{QGP} \simeq 1 \text{ fm}/c \sim t_{J/\psi}$
 - RHIC, LHC: $t_{QGP} \lesssim 0.1 \text{ fm}/c \sim t_{c\bar{c}}$

survival of initially-produced J/ψ at SPS/FAIR energies? ($T_d \sim T_c$)

- collision time, $t_{coll} = 2R/\gamma_{cm}$
 - SPS (FAIR): $t_{coll} \gtrsim t_{J/\psi}$
 - RHIC: $t_{coll} < t_{J/\psi}$, LHC: $t_{coll} \ll t_{J/\psi}$

cold nuclear suppression (breakup) important at SPS/FAIR energies?

shadowing is yet another (cold nuclear) effect - important at LHC (RHIC?)

NB: the only way to distinguish: measure $\sigma_{c\bar{c}}$ in pA and AA

Statistical hadronization: method and inputs

- Thermal model calculation (grand canonical) T, μ_B : $\rightarrow n_X^{th}$
- $N_{c\bar{c}}^{dir} = \frac{1}{2}g_c V (\sum_i n_{D_i}^{th} + n_{\Lambda_i}^{th}) + g_c^2 V (\sum_i n_{\psi_i}^{th} + n_{\chi_i}^{th})$
- $N_{c\bar{c}} \ll 1 \rightarrow$ Canonical (J.Cleymans, K.Redlich, E.Suhonen, Z. Phys. C51 (1991) 137):

$$N_{c\bar{c}}^{dir} = \frac{1}{2}g_c N_{oc}^{th} \frac{I_1(g_c N_{oc}^{th})}{I_0(g_c N_{oc}^{th})} + g_c^2 N_{c\bar{c}}^{th} \rightarrow g_c \text{ (charm fugacity)}$$

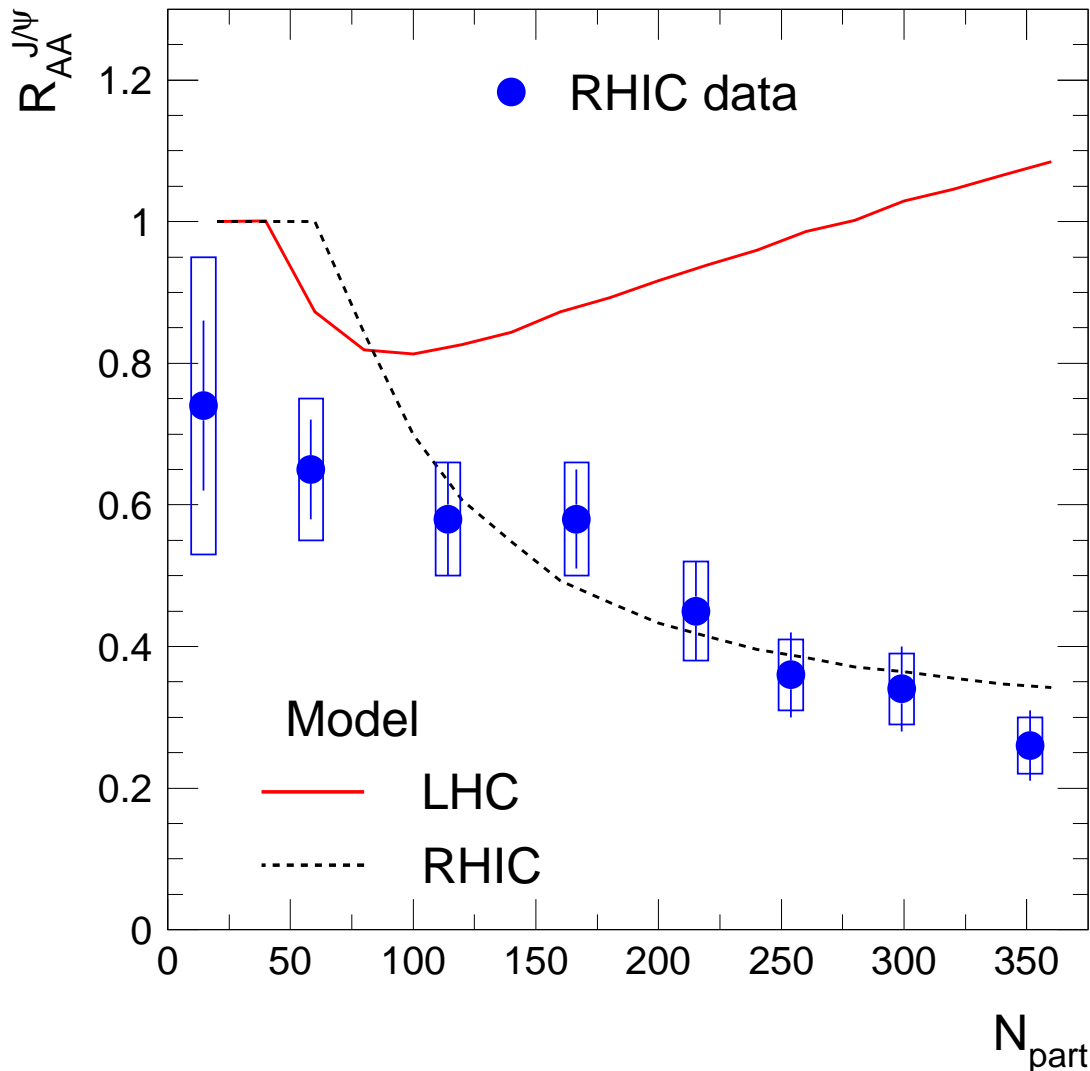
$$\text{Outcome: } N_D = g_c V n_D^{th} I_1/I_0 \quad N_{J/\psi} = g_c^2 V n_{J/\psi}^{th}$$

Inputs: $T, \mu_B, V_{\Delta y=1} (= (dN_{ch}^{exp}/dy)/n_{ch}^{th}), N_{c\bar{c}}^{dir}$ (pQCD or exp.)

Minimal volume for QGP: $V_{QGP}^{min} = 400 \text{ fm}^3$

Charmonium at RHIC and LHC

...an ultimate observable to measure the phase boundary (thermal model)
 ...with the help of charm quarks equilibrating in the deconfined stage



$$R_{AA}^{J/\psi} = (dN_{J/\psi}^{AuAu}/dy)/(N_{coll} \cdot dN_{J/\psi}^{pp}/dy)$$

$R_{AA}=1$ if superposition of pp coll.

very different centrality dependence

- "suppression" at RHIC

- "enhancement" at LHC

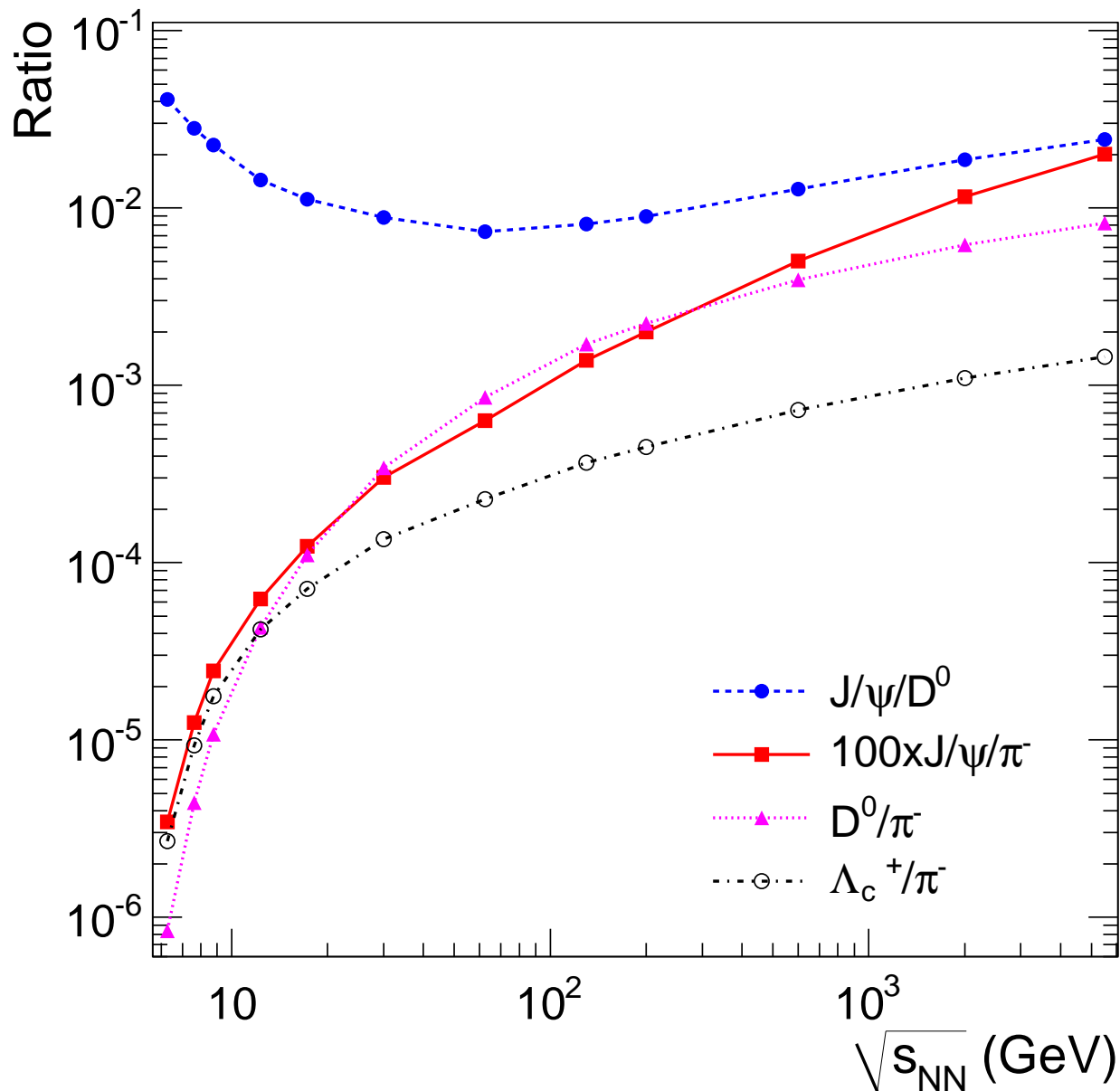
$$N_{J/\psi} \sim (N_{c\bar{c}}^{dir})^2$$

What is so different at LHC?

(compared to RHIC)

$\sigma_{c\bar{c}}$: 10x, Volume: 3x

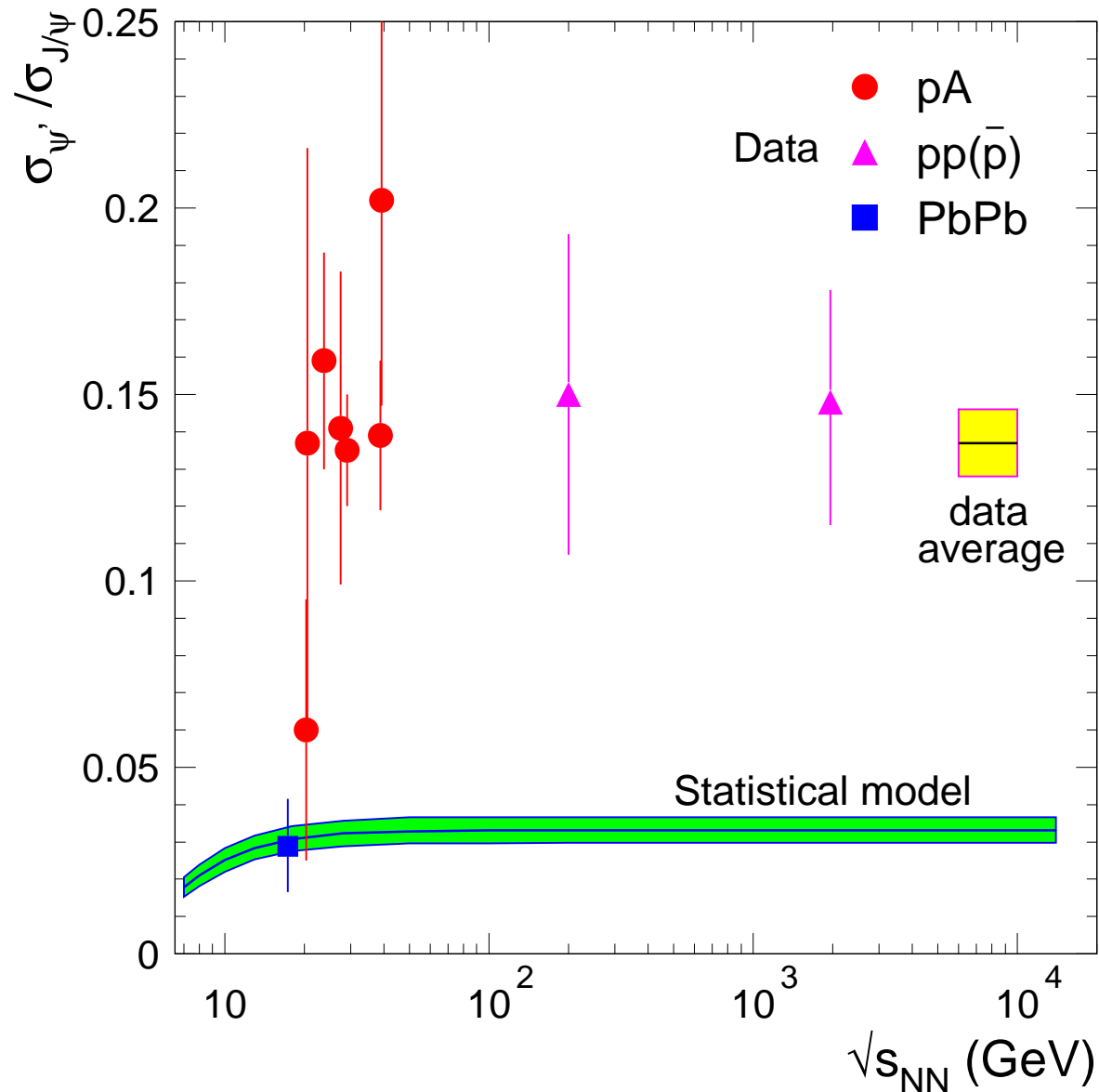
“Horns” for charmed hadrons?



...none ...

suppressed by the minute charm production at low energies

Charmonium in pp(A) collisions



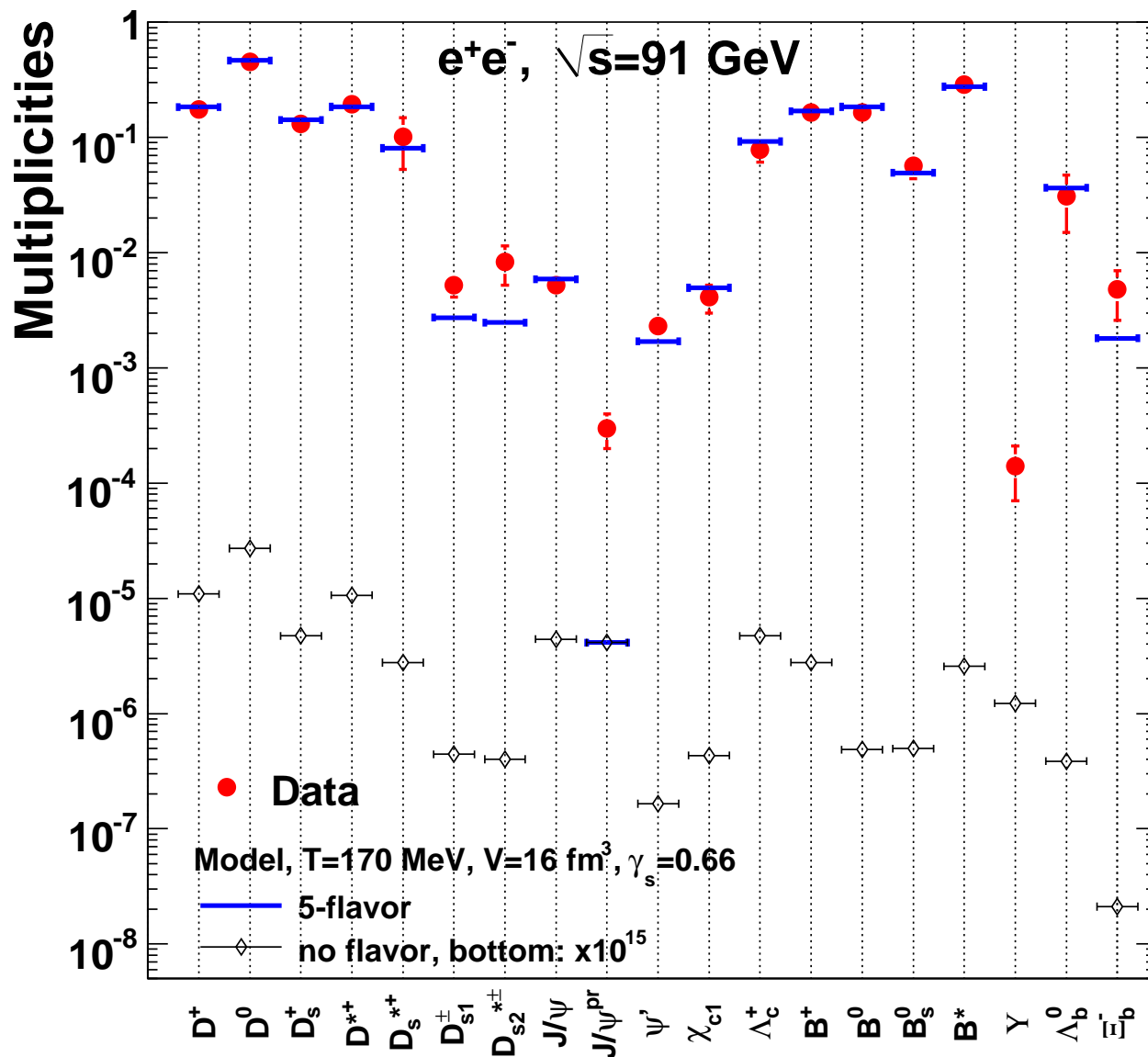
...is far from thermalized
(model is for AA)

...while a thermal value is
reached in central PbPb
(NA50, SPS)

more data in AA badly needed

NB: this ratio carries no sen-
sitivity to $\sigma_{c\bar{c}}$

Heavy quarks in e^+e^- collisions



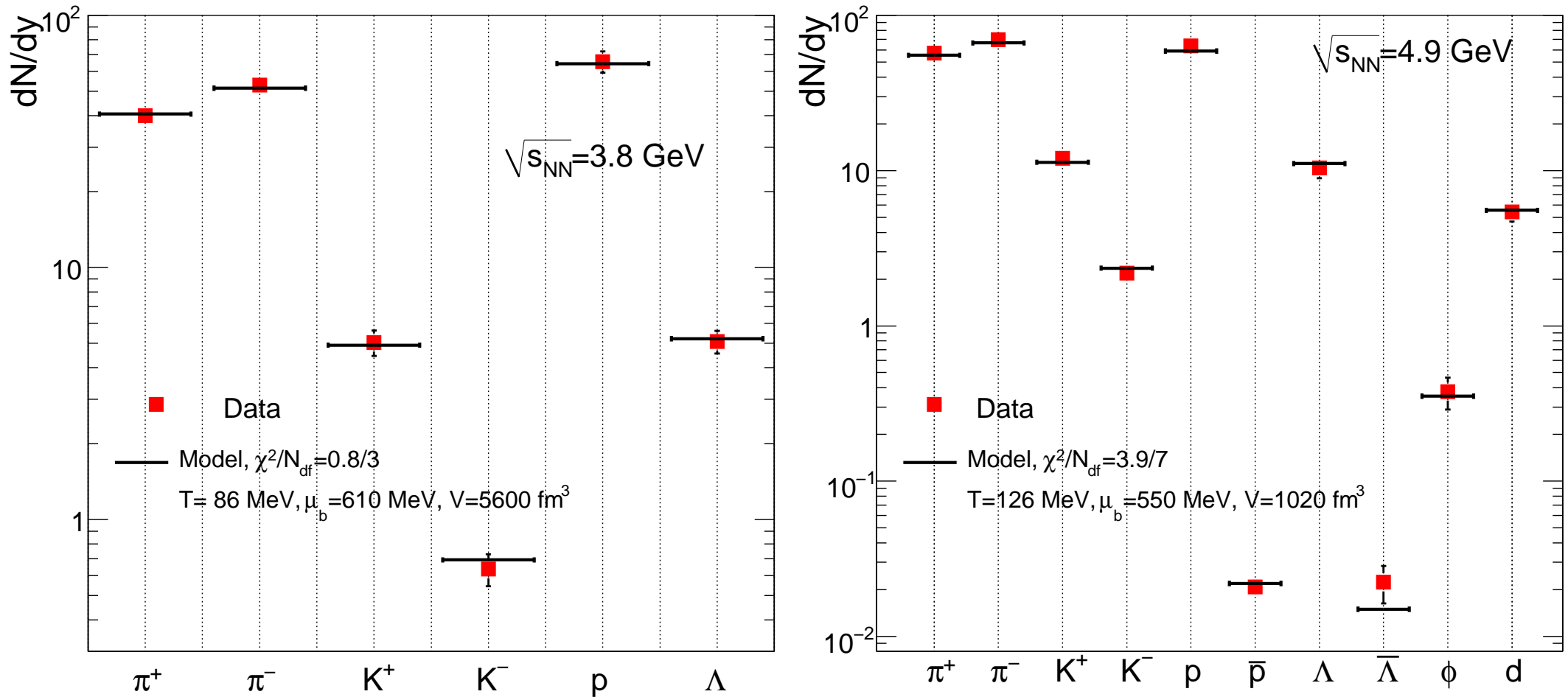
- open flavor hadrons strongly underpredicted in a pure thermal model (no flavor)
very different compared to u,d,s flavors
- agreement if $BR(Z^0 \rightarrow q\bar{q})$ are used in the model (5-flavor)!
(T, γ_s, V from fits of u,d,s flavors)
see also Becattini et al, EPJ C 56 (2008) 493
- quarkonia always strongly underpredicted

Summary

- thermal fits work remarkably well (AGS-RHIC) $\Rightarrow (T, \mu_b, V)$
- limiting temperature \Rightarrow phase boundary (LQCD)
 - \rightarrow for the skeptics... *LHC case will be decisive* ("bigger,...")
- indications (bad fits) around the critical point? ...maybe, at SPS...
 - ...but not a strong case due to disagreements between experiments
 - \rightarrow RHIC low-energy run (and CBM?) will clarify this
- Needed: a better freeze-out line (or phase boundary?) for $\mu_b > 500$ MeV
- no indications for strangeness non-equilibrium (γ_S) in central collisions (other models: not at SIS, RHIC; *some* at AGS-SPS, *some* at RHIC)
- *the model explains well charmonium (further tests soon at LHC)*
 - ...charm(onium) is then a powerful new observable to delineate the phase boundary

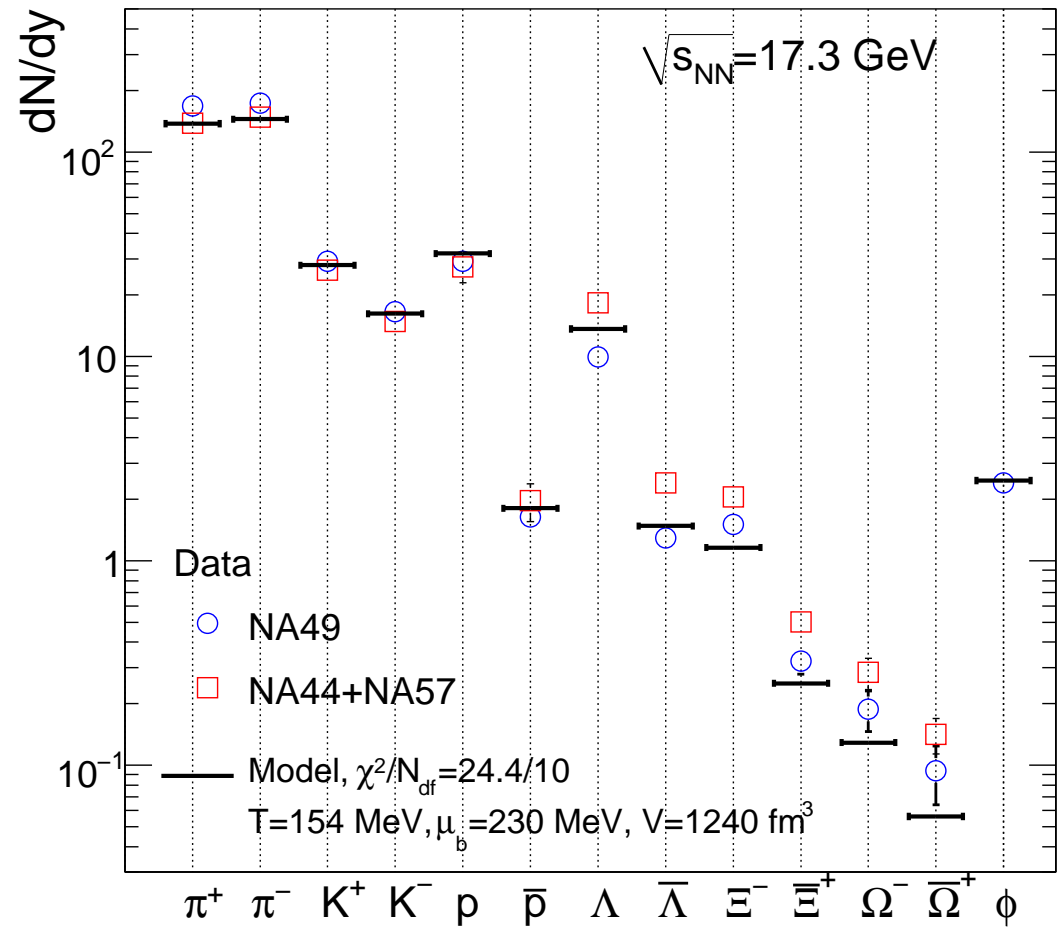
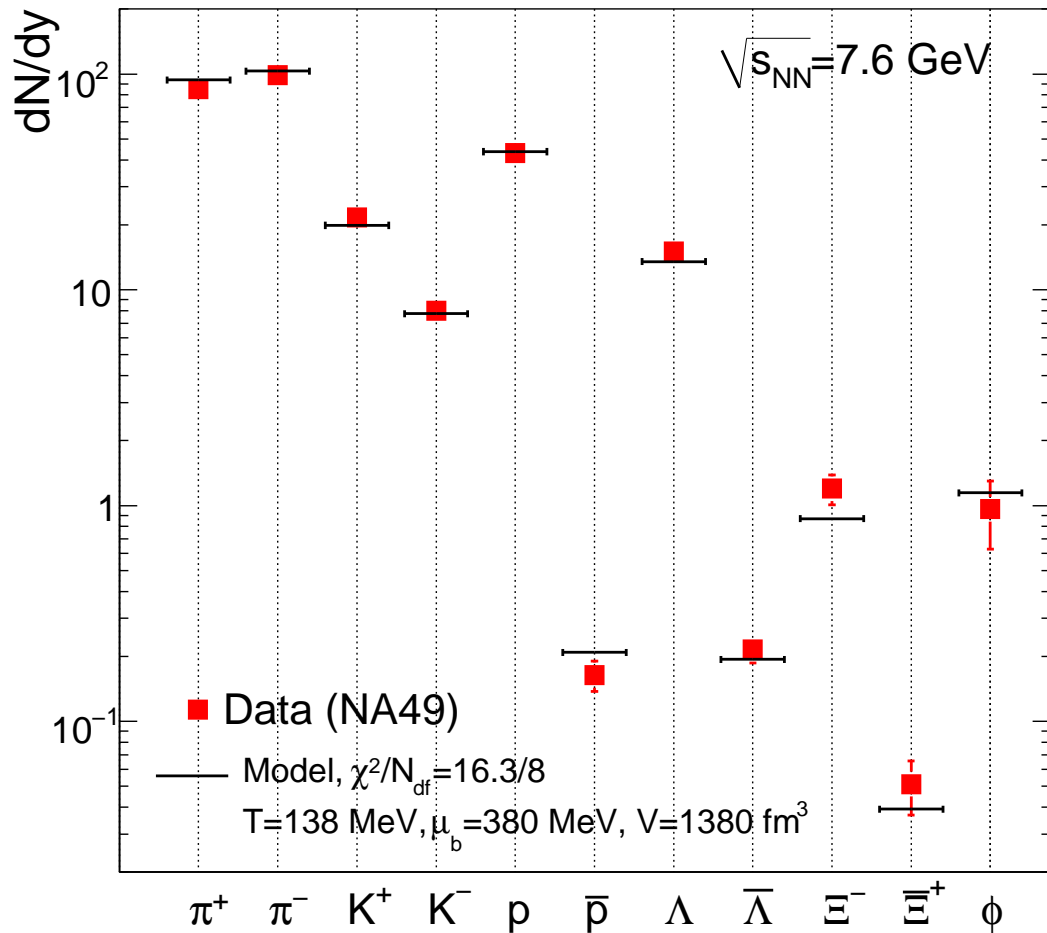
Backup slides

Fits at AGS: 6 and 10.5 AGeV



AGS, 2-8 AGeV: a rather small set of hadron yields measured

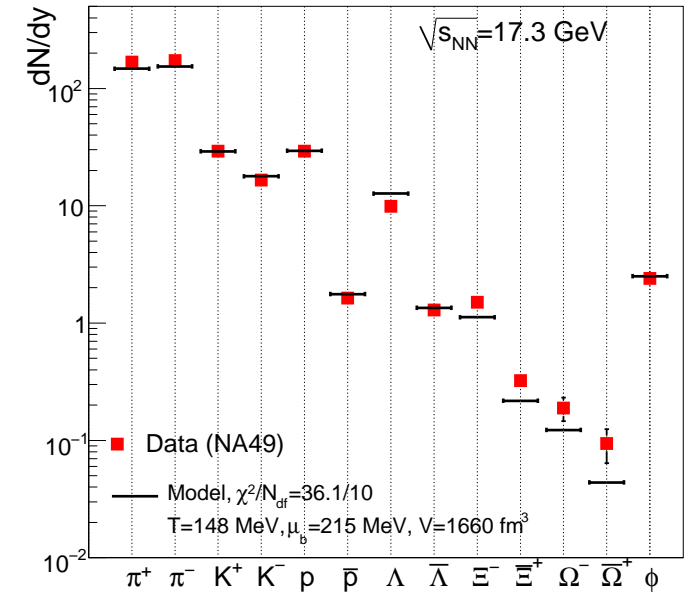
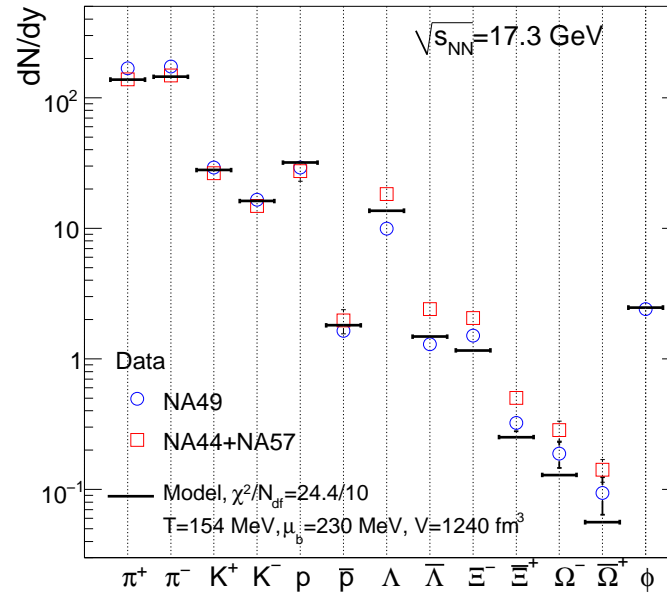
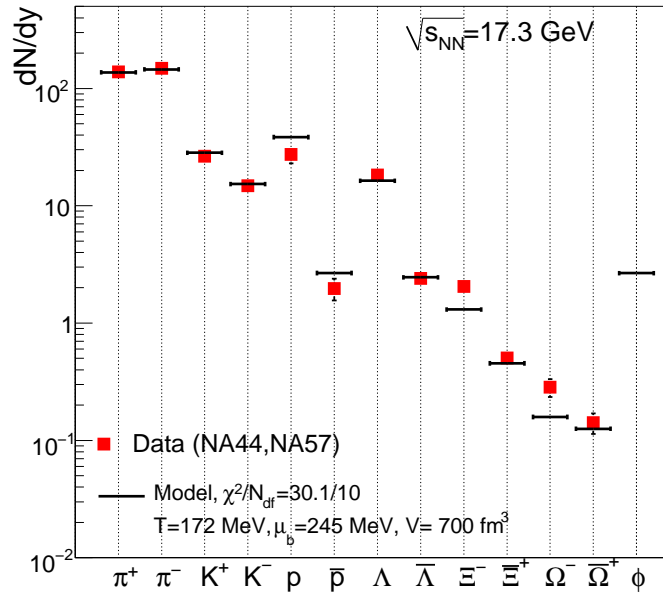
Fits at SPS: 30 and 158 GeV



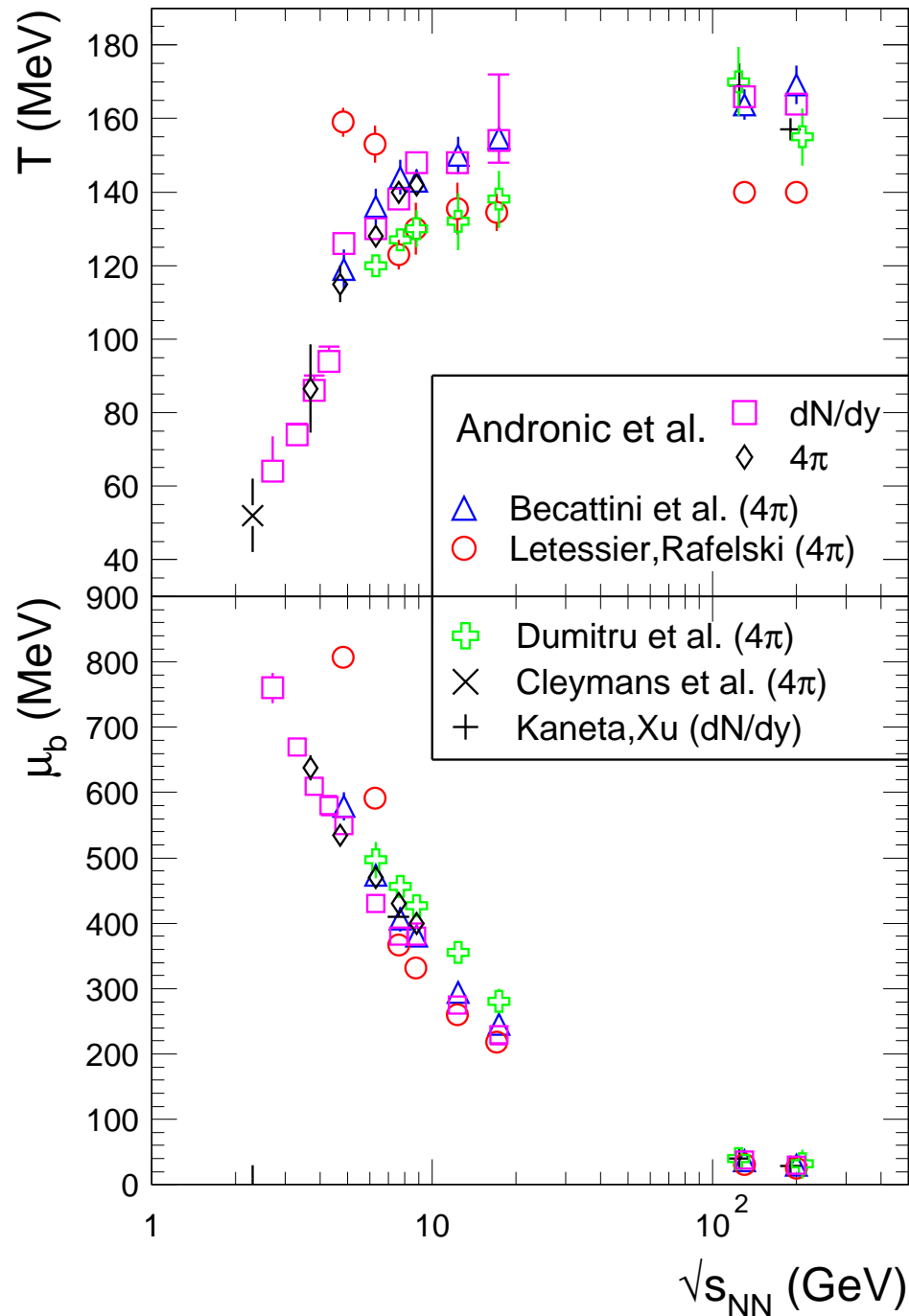
only NA49 data: $T = 148 \text{ MeV}, \mu_b = 215 \text{ MeV}, V = 1660 \text{ fm}^3, \chi^2/N_{df} = 36/10$

only NA44+NA57: $T = 172 \text{ MeV}, \mu_b = 245 \text{ MeV}, V = 700 \text{ fm}^3, \chi^2/N_{df} = 30/10$

SPS, 158 AGeV

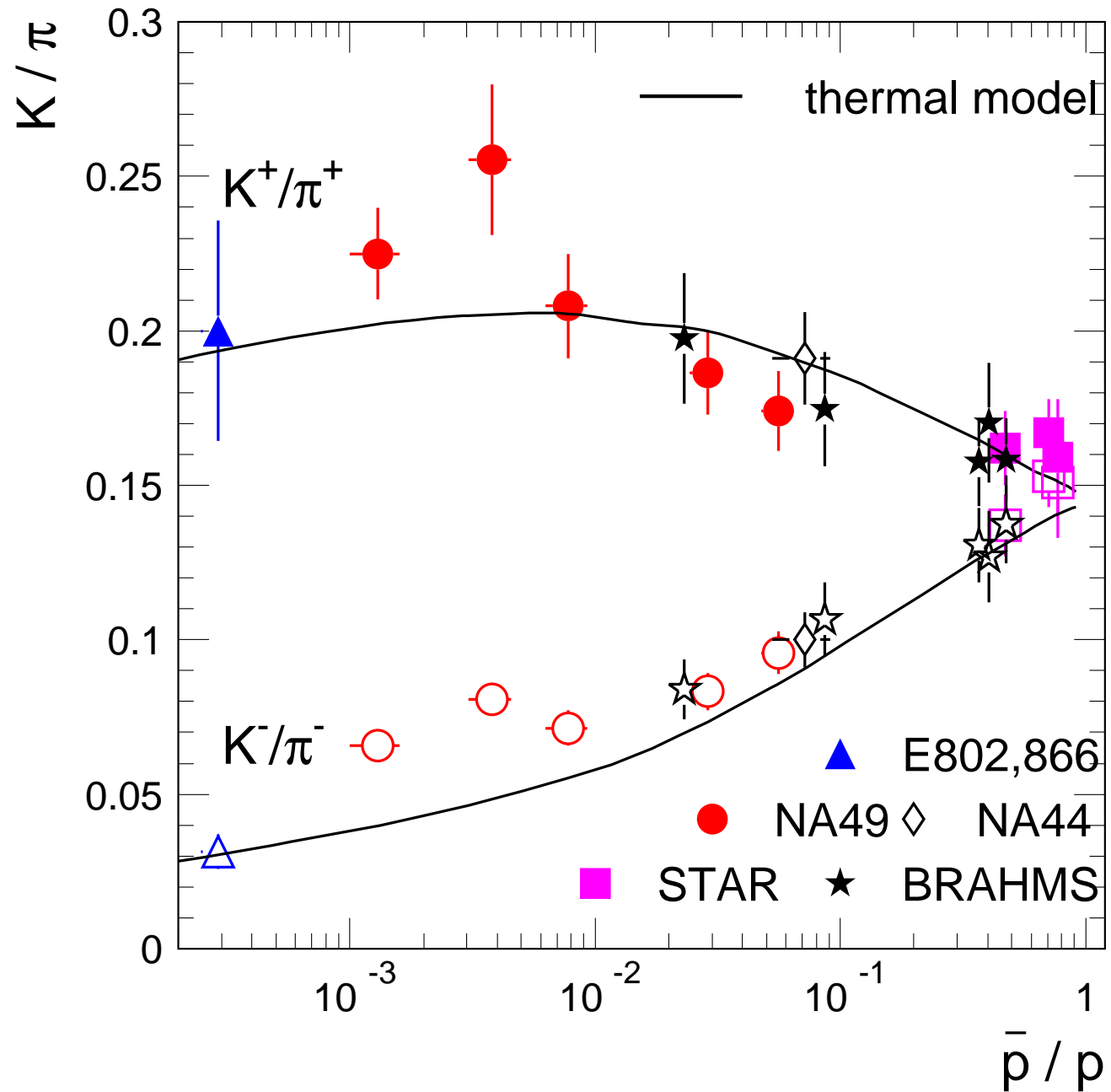


Energy dependence of the thermal parameters

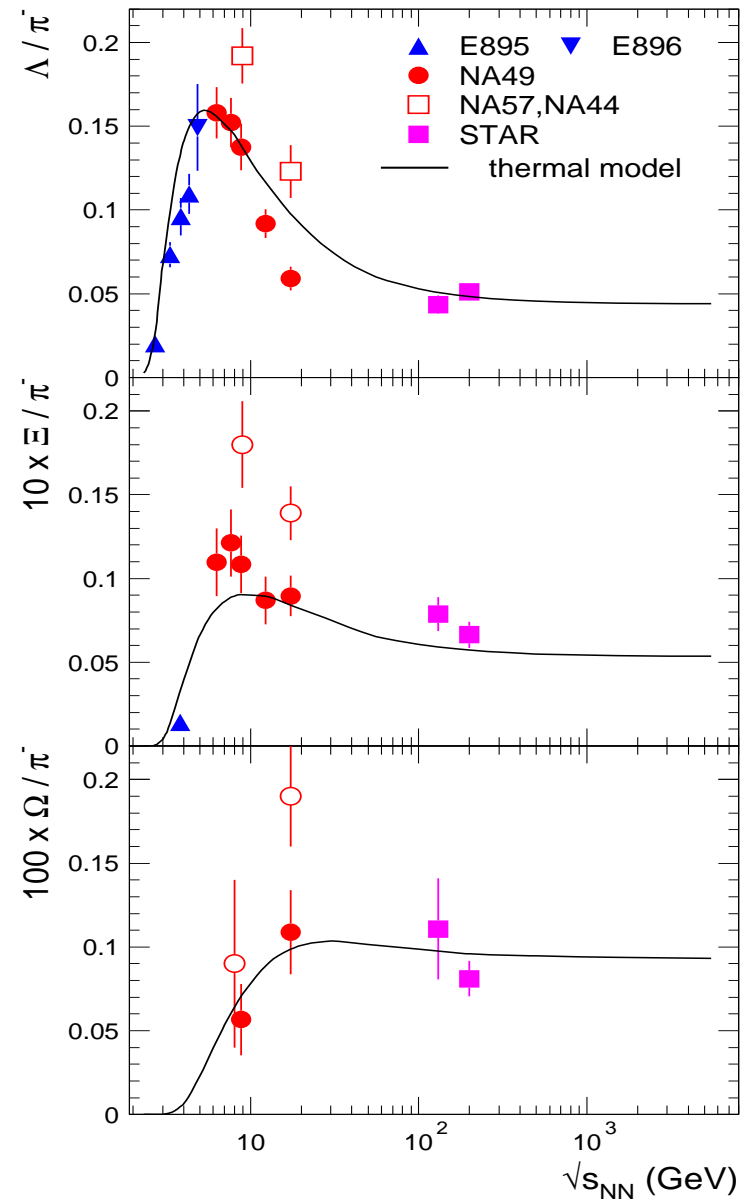
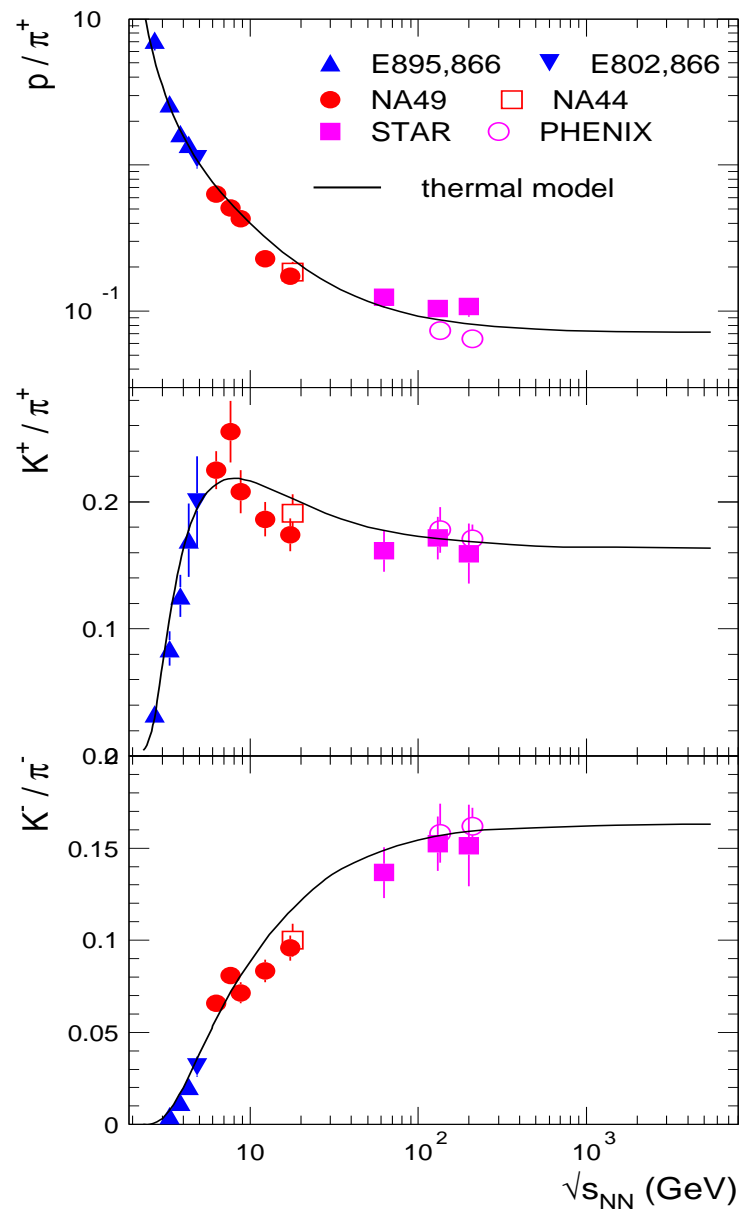


- Becattini et al.: $+\gamma_S$
Phys. Rev. C 73 (2006) 044905
Phys. Rev. C 78 (2008) 054901
- Rafelski et al.: $+\gamma_{S,q}, \lambda_{q,S,I_3}$
Eur. Phys. J. A 35 (2008) 221
 $\gamma_S=0.18, 0.36, 1.72, 1.64, \dots$
 $\gamma_q=0.33, 0.48, 1.74, 1.49, 1.39, 1.47, \dots$
- Dumitru et al.: inhomogeneous freeze-out
($\delta T, \delta \mu_B$)
Phys. Rev. C 73 (2006) 024902
- Kaneta, Xu, nucl-th/0405068
- Cleymans et al., Phys. Rev. C 57 (1998) 3319

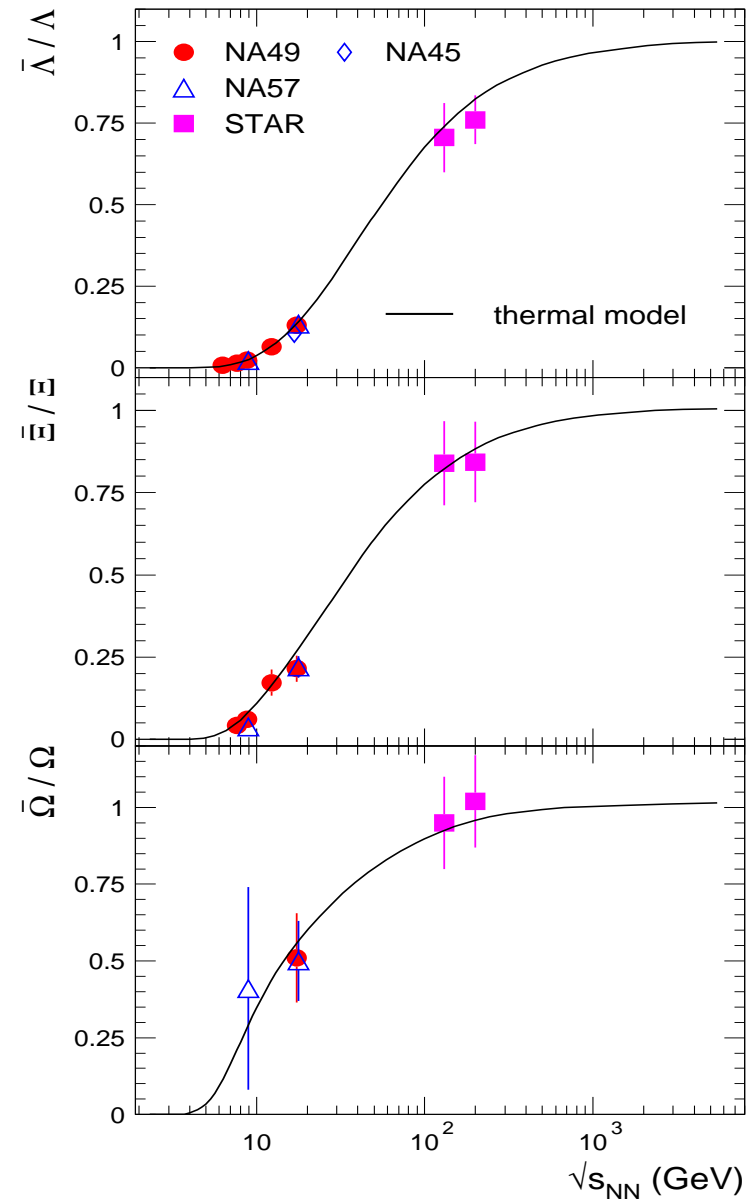
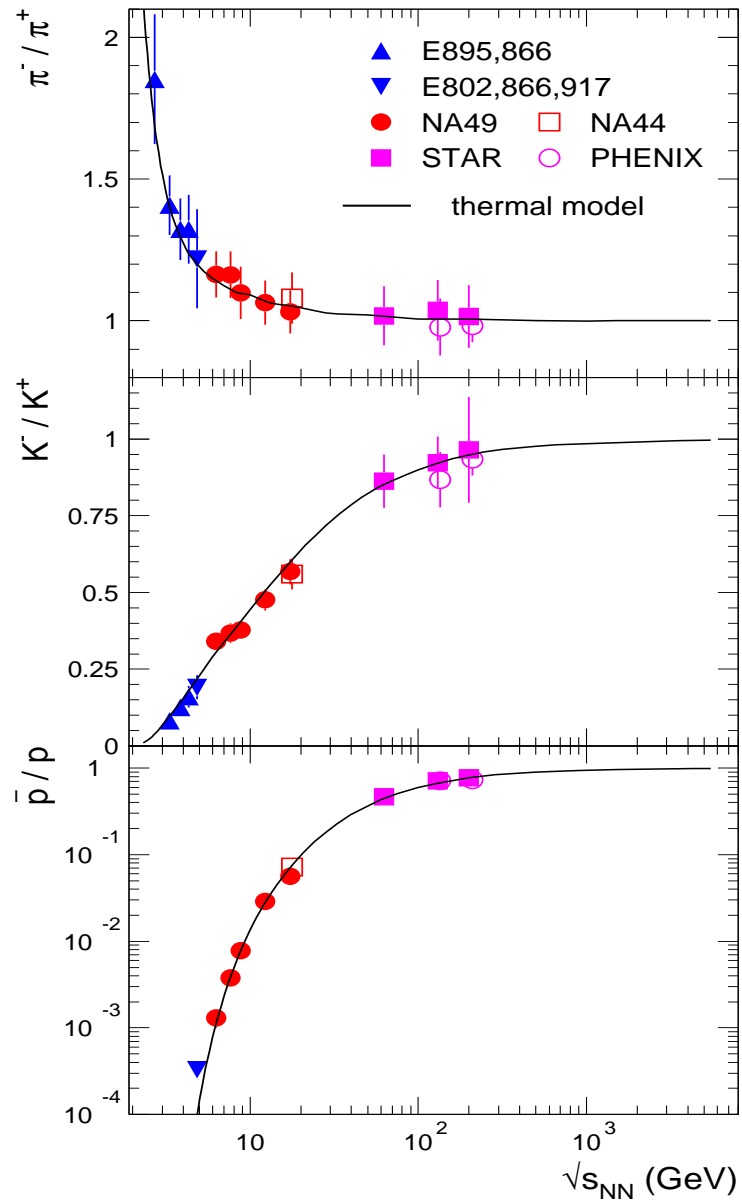
A clever proxy for the horn (BRAHMS collab.)



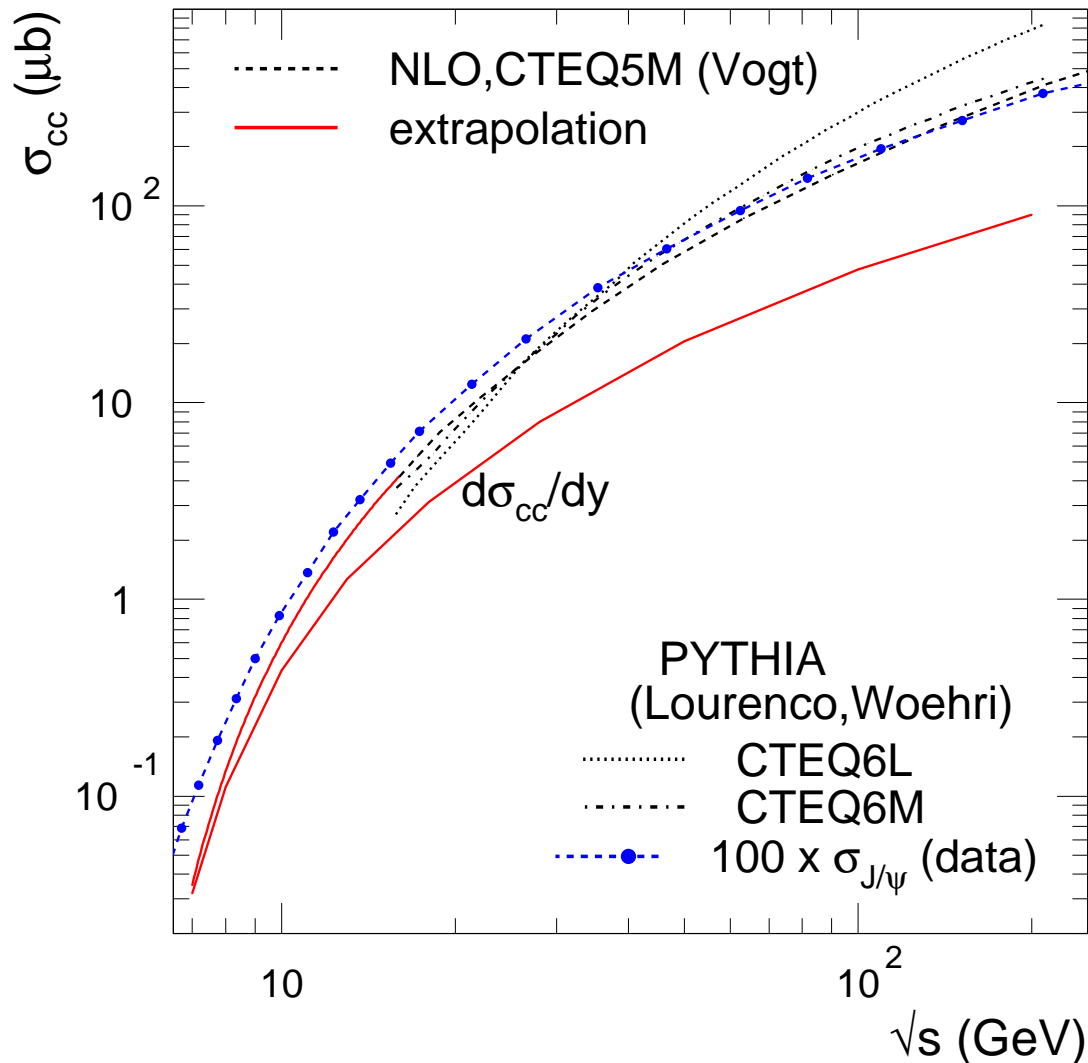
More particle ratios



More particle ratios



$N_{c\bar{c}}^{dir}$ from pQCD calculations (pp)



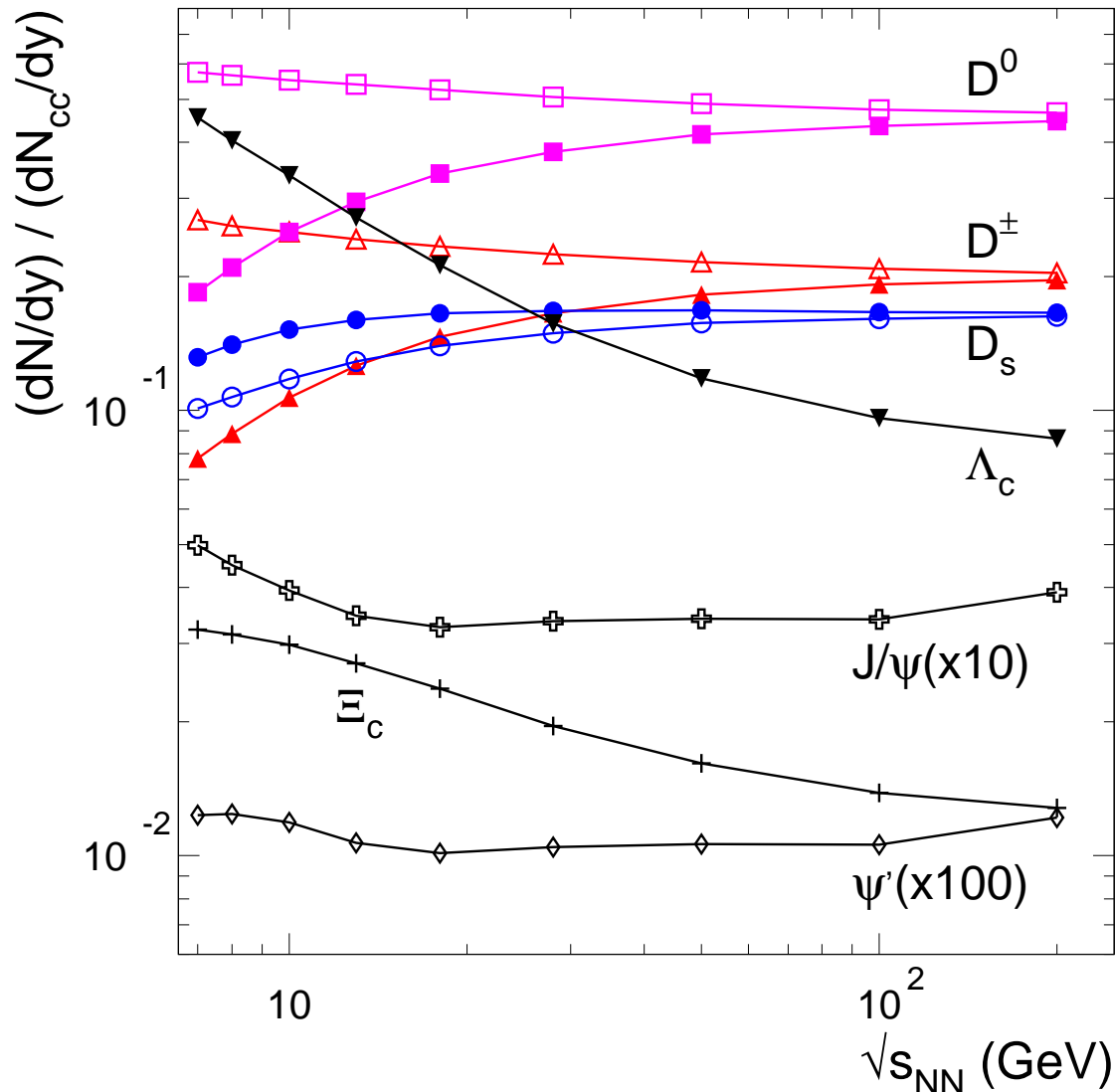
R.Vogt, IJMP E12 (2003) 211
[hep-ph/0111271]

pQCD is not parameter-free!
(PDF, m_c , μ_R , μ_F)

$dN_{c\bar{c}}/dy$ for central collisions
($N_{part}=350$):

RHIC: $\simeq 1.6$, LHC: $\simeq 16$

Overall charm chemistry



yields per initial charm pair

- Λ_c prod. favored at large μ_b
...it's a must at FAIR (CBM)
- isospin is important
- ψ'/ψ relative yield:
3% in QGP, 13% in pp
decreases at low energies
 $\sqrt{s_{NN}}=7-10$ GeV:
 $T=151-161$ MeV
- charmed hadrons can signal the onset of QGP