

Photon and hadron observables from a 3+1 dimensional hydro model



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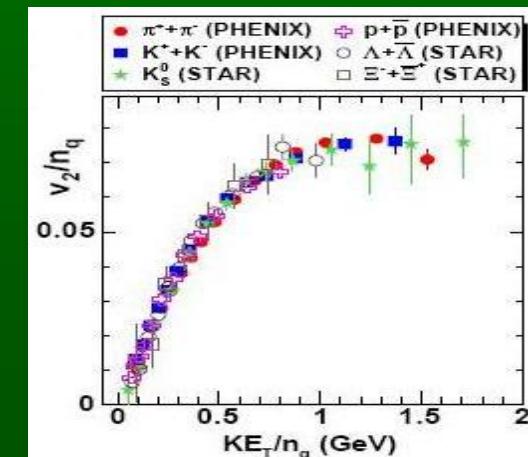
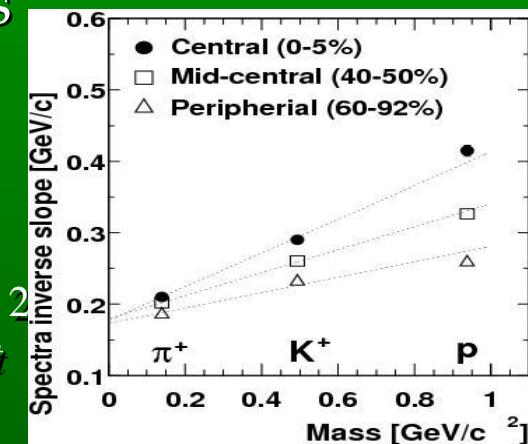
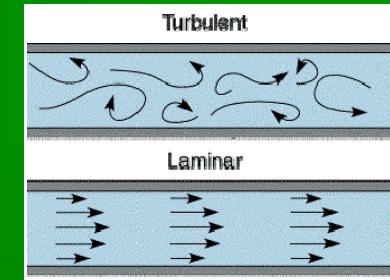
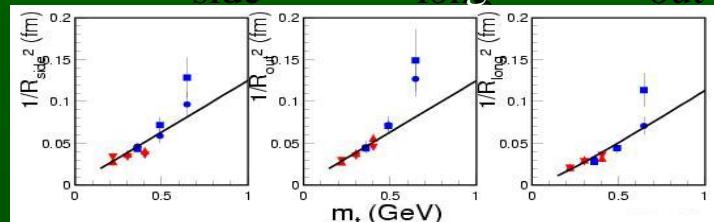
- A 3+1 dimensional model
- Hadron spectra, flow and HBT
- Photon observables

VI Workshop on Particle Correlations and Femtoscopy

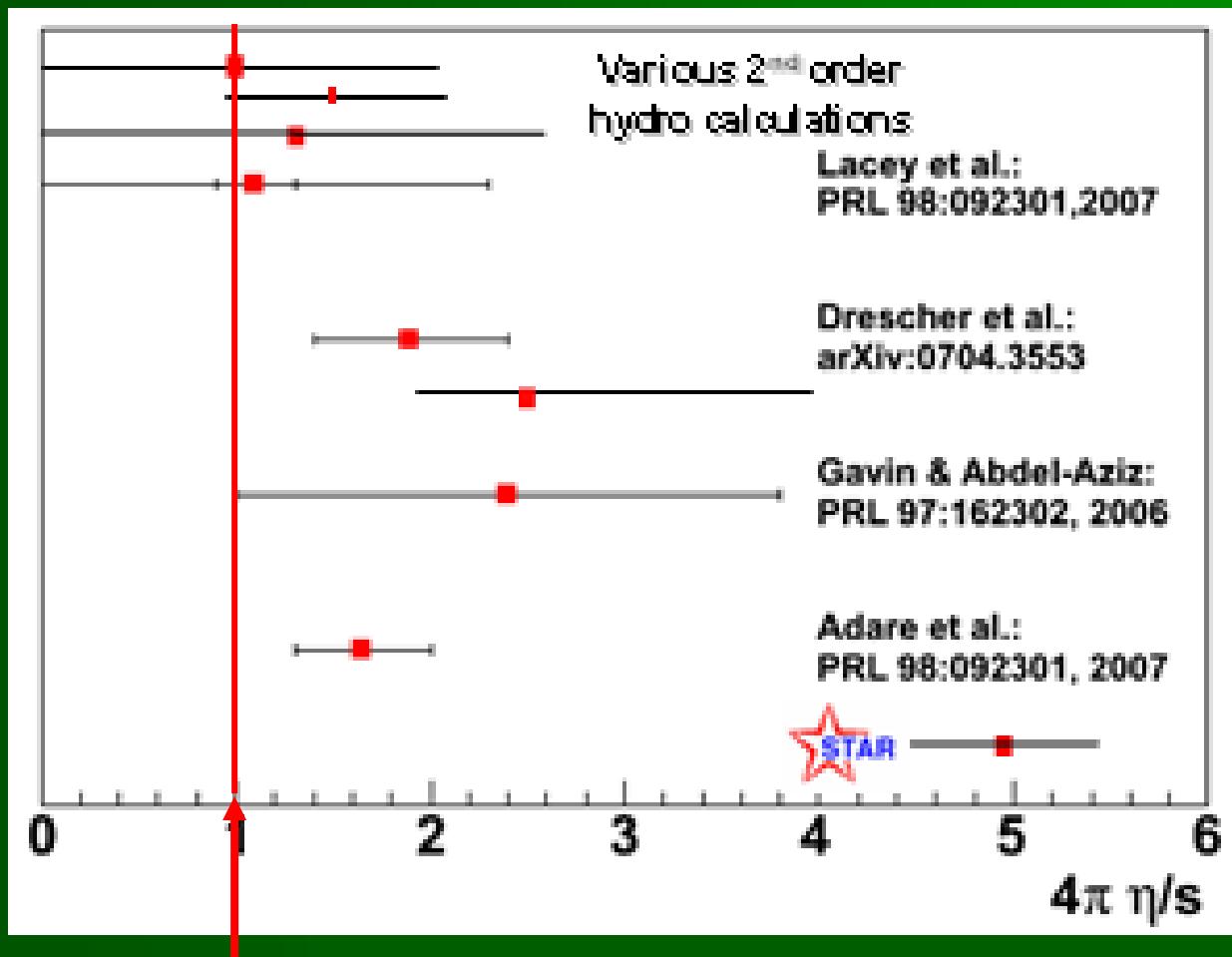
Kiev, September 14-18, 2010

Hydrodynamic scaling predictions

- Hydro predicts scaling (also viscous)
- What does a scaling mean?
 - For example Reynolds number $\rho v r / \eta$
 - Only a combination of parameters matters
- Collective, thermal behavior →
Loss of information
- Spectra slopes: $N_1 \sim e^{-\frac{m_t}{T_{\text{eff}}}}$; $T_{\text{eff}} = T_0 + m u_t^2$
- Elliptic flow: $v_2 = \frac{I_1(w)}{I_0(w)}$; $w \sim KE_T$
- HBT radii: $R_{\text{side}}^2 \approx R_{\text{long}}^2 \approx R_{\text{out}}^2 \sim \frac{1}{m_t}$



Perfect hydro picture

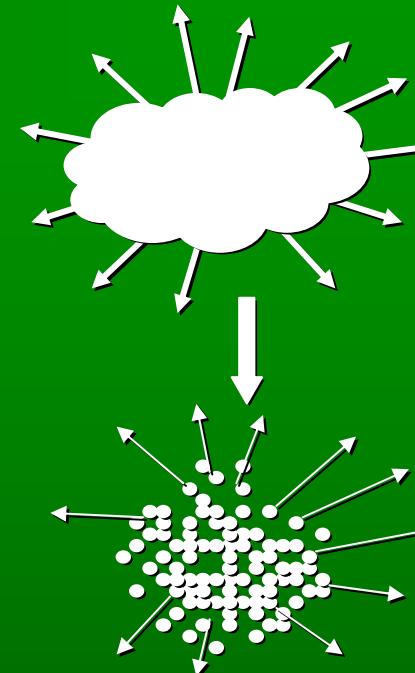


- No data point even near the kinematic viscosity of ${}^4\text{He}$ ($10/(4\pi)$)
- Close to AdS/CFT minimum, $(1/(4\pi))$

Little vocabulary of hydrodynamics

- Exact/parametric solution
 - Solution of hydro equations analytically, without approximation
 - Usually has free parameters
- Hydro inspired parameterization
 - Distribution determined at freeze-out only, their time dependence is not considered
- Numerical solution
 - Solution of hydro equations numerically
 - Start from arbitrary initial state

How analytic hydro works



- Take hydro equations and EoS
- Find a solution
 - Will contain parameters (like Friedmann, Schwarzschild etc.)
 - Will use a possible set of initial conditions
- Use a freeze-out condition
 - Eg fixed proper time or fixed temperature
 - Generally a hyper-surface
- Calculate the hadron source function
- Integrate over time to calculate photon spectra
- Hydrodynamics: Initial conditions \otimes dynamical equations \otimes freeze-out conditions

Some known relativistic solutions

Solution	Basic prop's	EoS	Observables
Csörgő, Nagy, Csanád Phys.Lett.B 663:306-311, 2008 Phys.Rev.C77:024908,2008	Ellipsoidal, 1D accelerating	$\varepsilon - B = \kappa(p + B)$	$dn/dy, \varepsilon$
Landau Izv. Acad. Nauk SSSR 81 (1953) 51	Cylindr., 1D, accelerating	$\varepsilon = \kappa n T$	none
Hwa-Björken R.C. Hwa, PRD10, 2260,1974 J.D. Bjorken, PRD27, 40(1983)	Cylindr., 1D, non-accelerating	$\varepsilon = \kappa n T$	$dn/dy, \varepsilon$
Bialas et al. A. Bialas, R. A. Janik, and R. B. Peschanski, Phys. Rev. C76, 054901 (2007).	1D, between Landau and Hwa-Björken	$\varepsilon = \kappa n T$	dn/dy
Csörgő, Csernai, Hama, Kodama Heavy Ion Phys. A 21, 73 (2004))	Ellipsoidal, 3D, non-accelerating	$\varepsilon = \kappa n T$	This work does the calculation

Where we are

- Revival of interest, new solutions
 - Sinyukov, Karpenko, nucl-th/0505041
 - Pratt, nucl-th/0612010
 - Bialas et al.: Phys.Rev.C76:054901,2007
 - Borsch, Zhdanov: SIGMA 3:116,2007
 - Nagy et al.: J.Phys.G35:104128,2008 and arXiv/0909.4285
 - Liao et al.: arXiv/09092284 and Phys.Rev.C80:034904,2009
 - Mizoguchi et al.: Eur.Phys.J.A40:99-108,2009
 - Beuf et al.:Phys.Rev.C78:064909,2008 (dS/dy as well!)
- Need for solutions that are:
 - accelerating + relativistic+ 3 dimensional
 - explicit + simple + compatible with the data
- Need to calculate observables!

The solution we investigate

- Density, temperature, pressure

- $v(s)$ arbitrary, but realistic to choose Gaussian $v(s) = e^{-bs/2}$
 $b < 0$ is realistic

$$n = n_0 \left(\frac{\tau_0}{\tau} \right)^3 v(s)$$

$$T = T_0 \left(\frac{\tau_0}{\tau} \right)^{(3/\kappa)} \frac{1}{v(s)}$$

$$p = p_0 \left(\frac{\tau_0}{\tau} \right)^{\left(3 + \frac{3}{\kappa} \right)}$$

$$u^\mu = \gamma \left(1, \frac{\dot{X}(t)}{X(t)} x, \frac{\dot{Y}(t)}{Y(t)} y, \frac{\dot{Z}(t)}{Z(t)} z \right)$$

- Ellipsoidal symmetry

$$s = \frac{x^2}{X(t)^2} + \frac{y^2}{Y(t)^2} + \frac{z^2}{Z(t)^2}$$

- (thermodynamic quantities const. on the $s=\text{const.}$ ellipsoid)

- Directional Hubble-flow

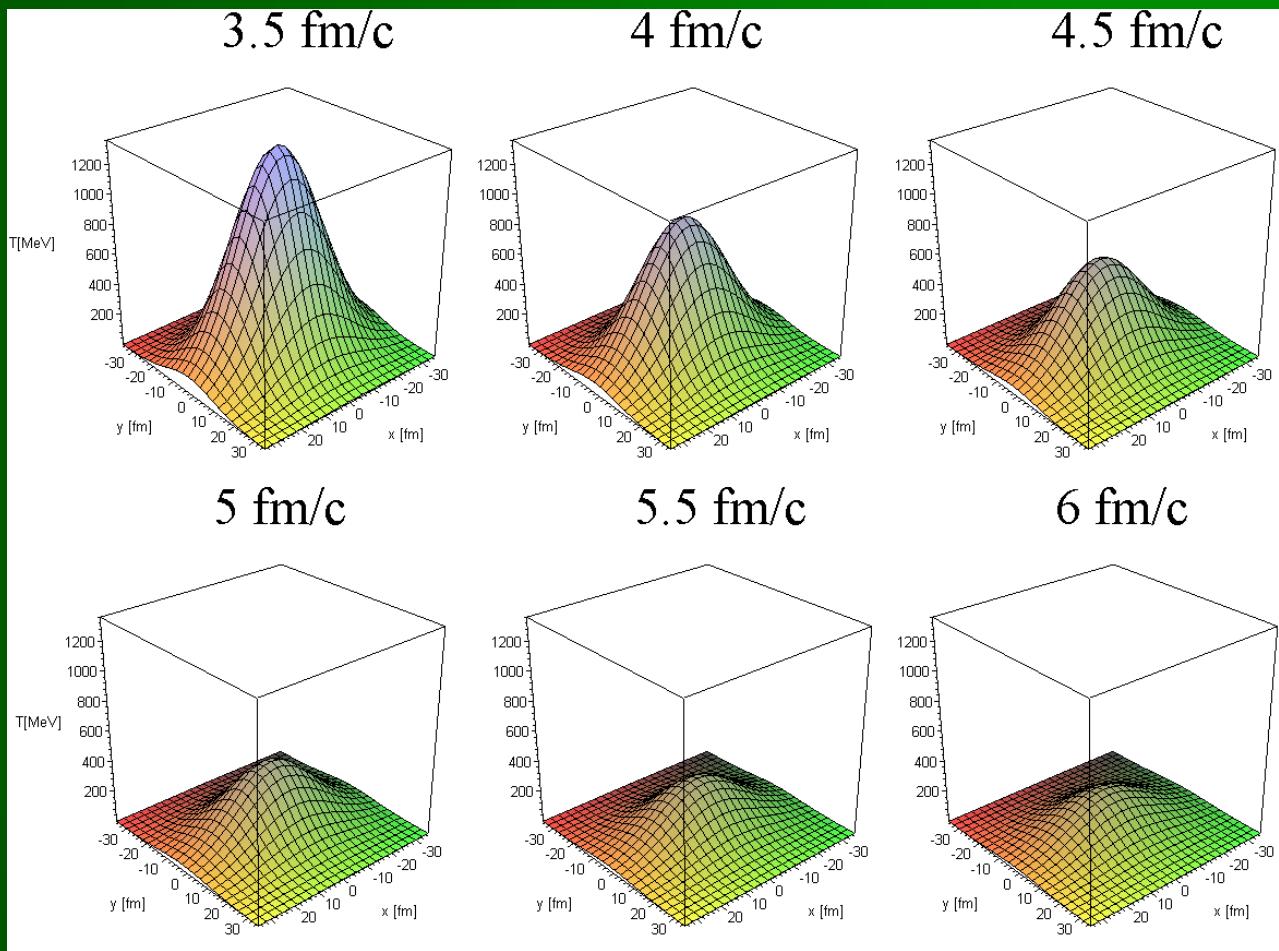
- $v = Hr$ or $H = v/r$, the Hubble-constants:
 - $\dot{X}(t), \dot{Y}(t), \dot{Z}(t) : \text{const.}; (\dot{X}, \dot{Y}) \Leftrightarrow (u_t, \varepsilon)$

$$\frac{\dot{X}(t)}{X(t)}, \frac{\dot{Y}(t)}{Y(t)}, \frac{\dot{Z}(t)}{Z(t)}$$

(T. Csörgő, L. P. Csernai, Y. Hama és T. Kodama, Heavy Ion Phys. A 21, 73 (2004))

Temperature profile

Transverse temperature profile as a function of time
with an example parameter set:



The hadronic source function

- Source function: probability of a particle created at x with momentum p
- Maxwell-Boltzmann distribution + extra terms

$$S(x, p)d^4x \sim n(x) \exp\left[-\frac{p_\mu u^\mu(x)}{T(x)}\right] \frac{p_\mu u^\mu}{u^0} H(\tau) d^4x$$

- $H(\tau)d\tau$ freeze-out distribution
if sudden: $H(\tau) = \delta(\tau - \tau_0)$
- $\frac{p_\mu u^\mu}{u^0} d^3x$ Cooper-Fry prefactor (flux term)
- Validity: $\tau_0 > R_{\text{HBT}}$, $m_t > T_0$

Hadronic results

- Single particle spectrum

$$N_1(p_t) = \bar{N} \bar{V} \left(m_t - \frac{p_t^2(T_{\text{eff}} - T_0)}{m_t T_{\text{eff}}} \right) \exp \left[-\frac{p_t^2}{2m_t T_{\text{eff}}} + \frac{p_t^2}{2m_t T_0} - \frac{m_t}{T_0} \right]$$

$$T_x = T_0 + m_t \dot{X}^2 \frac{T_0}{b(T_0 - E)}, \quad T_y = T_0 + m_t \dot{Y}^2 \frac{T_0}{b(T_0 - E)}, \quad \frac{1}{T_{\text{eff}}} = \frac{1}{2} \left(\frac{1}{T_x} + \frac{1}{T_y} \right)$$

- Elliptic flow

$$v_2 = \frac{I_1(w)}{I_0(w)} \quad \text{with} \quad w = \frac{p_t^2}{4m_t} \left(\frac{1}{T_y} - \frac{1}{T_x} \right) \sim E_K \frac{\varepsilon}{T_{\text{eff}}}, \quad I: \text{Bessel func.}$$

- Two-particle correlation functions: $R_{x,y} = \frac{T_0 \tau_0 (T_{x,y} - T_0)}{m_t T_{x,y}}$

$$R_{\text{out}} = R_{\text{side}} = 0.5 \left(R_x^2 + R_y^2 \right)$$

Source function of photons

- Photons are continuously created from thermalization
- Photons are not thermalized
- Bose-Einstein distribution

$$S(x, p)d^4x \sim \left(e^{p_\mu u^\mu / T} - 1 \right)^{-1} E d^4x$$

- Has to be integrated over time
- Second order Gaussian approximation

Photon spectra and photon v_2

- Integration can be done analytically

$$N_1(p_t) = (2\pi)^{3/2} p_t R_x R_y R_z \tau_0 \frac{\kappa}{3} A^{-4\kappa/3} \left[A^{3/2} \Gamma\left(-\frac{3}{2} + \frac{4\kappa}{3}, A\xi^{\frac{3}{\kappa}}\right) \Big|_{1}^{\tau_f/\tau_i} \right. \\ \left. + \frac{B^2}{4} A^{-1/2} \Gamma\left(\frac{1}{2} + \frac{4\kappa}{3}, A\xi^{\frac{3}{\kappa}}\right) \Big|_{1}^{\tau_f/\tau_i} \right]$$

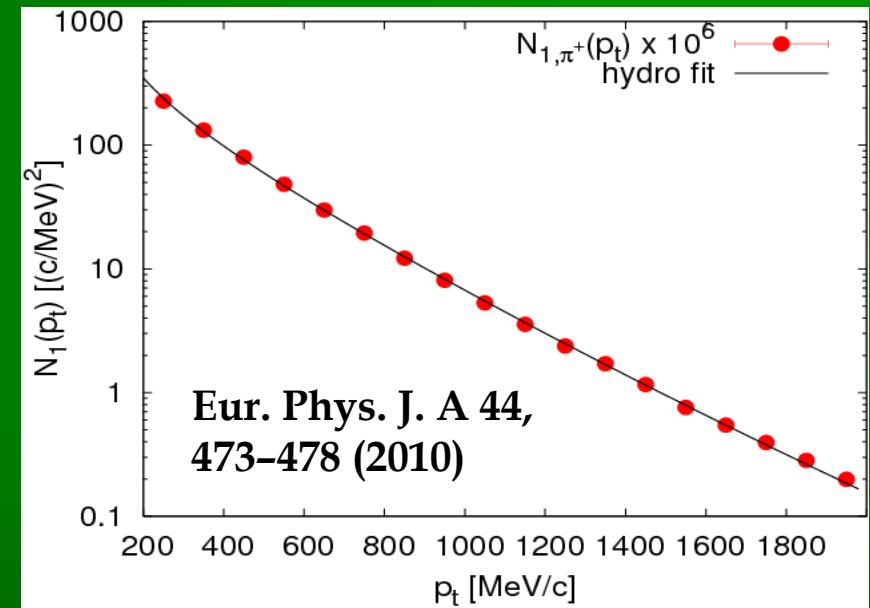
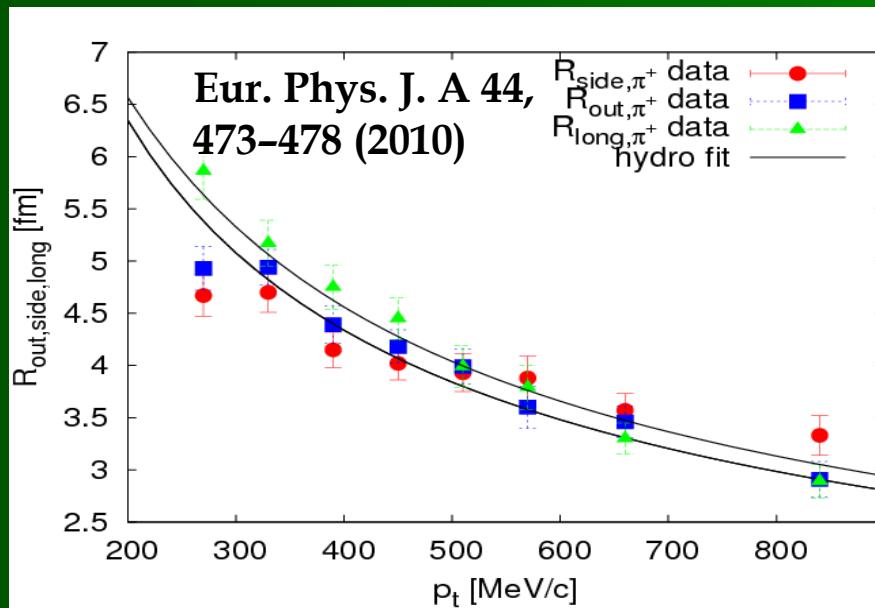
$$v_2 = \frac{B}{2A^2} \cdot \frac{\Gamma\left(-\frac{1}{2} + \frac{4\kappa}{3}, A\xi^{\frac{3}{\kappa}}\right) \Big|_{1}^{\tau_f/\tau_i}}{\Gamma\left(\frac{3}{2} + \frac{4\kappa}{3}, A\xi^{\frac{3}{\kappa}}\right) \Big|_{1}^{\tau_f/\tau_i}}$$

- A and B are:

$$A = \frac{p_t}{T_0} \left[1 - \frac{\kappa(\kappa-3) - \frac{\kappa^2 b}{u_t^2}}{2 \left((\kappa-3) - \frac{b}{u_t^2} \right)} \right], \quad B = \frac{p_t}{T_0} \frac{\kappa^2 b \frac{\varepsilon}{u_t^2}}{2 \left((\kappa-3) - \frac{b}{u_t^2} \right)}$$

Single pion spectrum with HBT radii

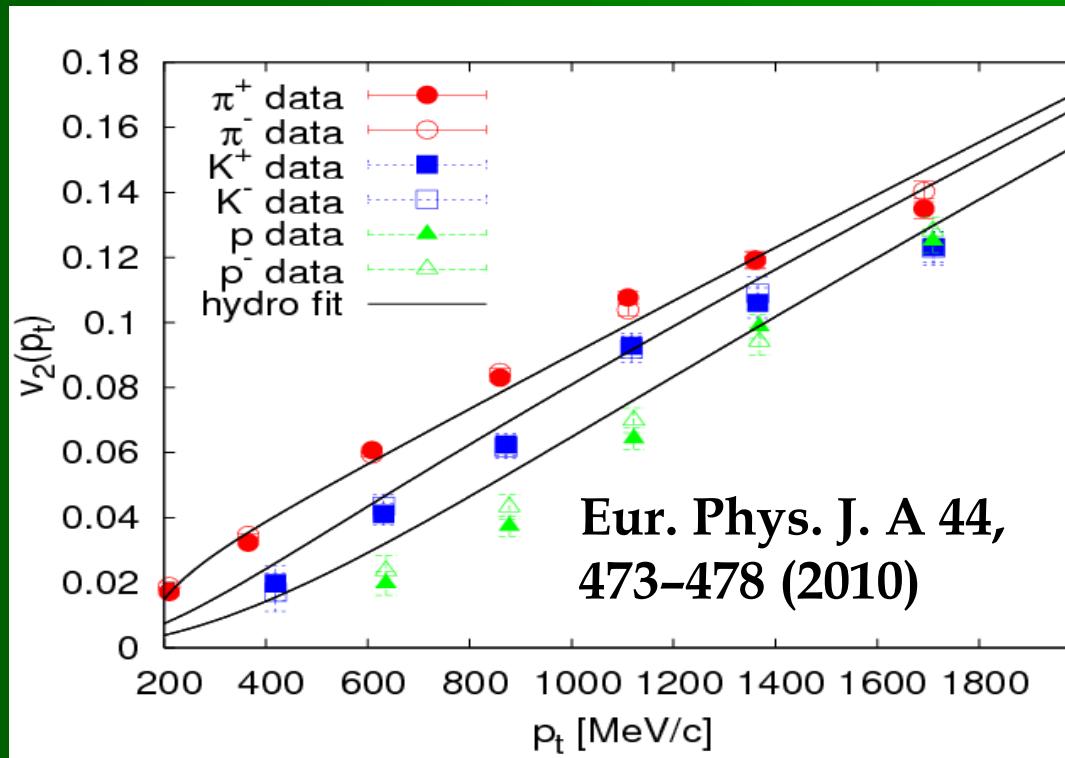
- 0-30% centrality, Au+Au, PHENIX



- T_0 199 ± 3 MeV central freeze-out temp.
- ε 0.80 ± 0.02 momentum space ecc.
- u_t^2/b -0.84 ± 0.1 ($b < 0$) transv. flow / temp. grad
- τ_0 7.7 ± 0.1 freeze-out proper time
- χ^2 171 (24 with theory error)

Elliptic flow

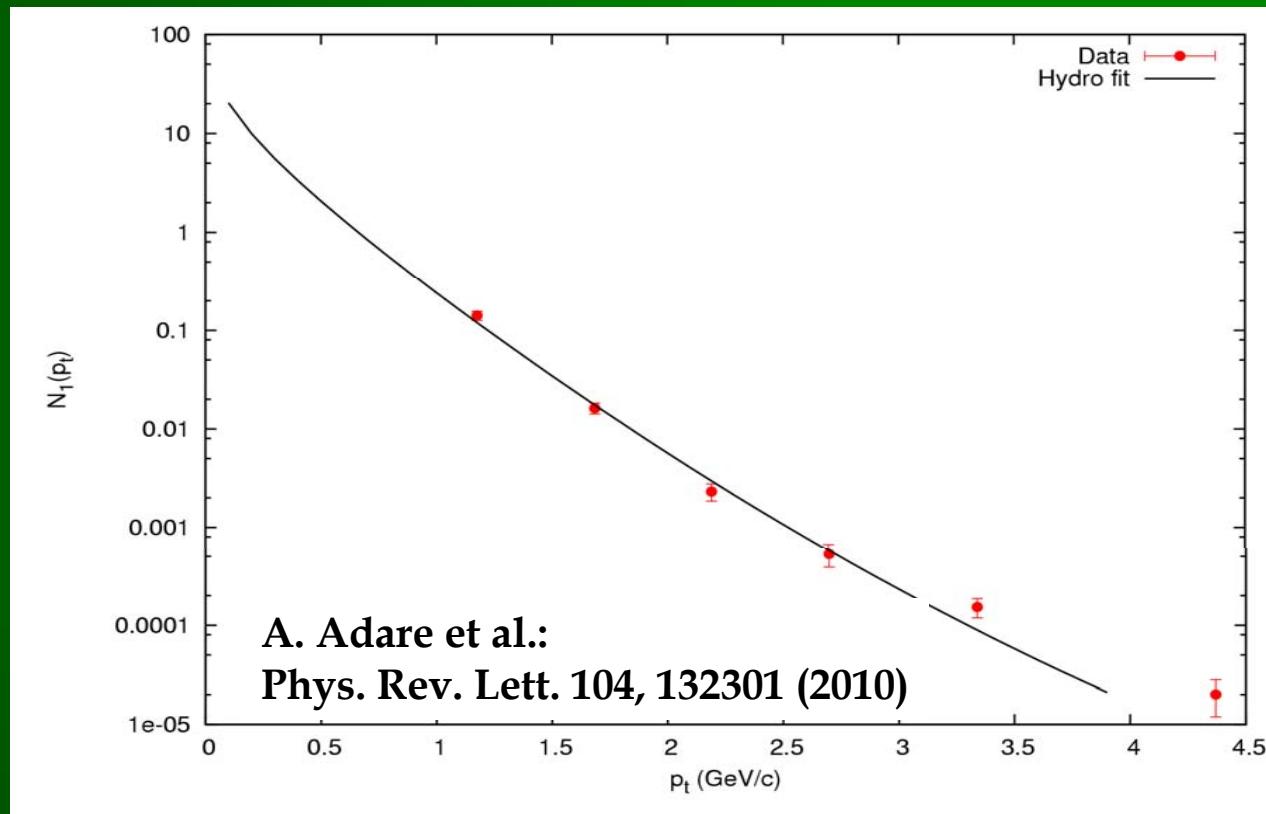
- 0-92% centrality, Au+Au, PHENIX (R.P. method technique)



- T_0 204 ± 7 MeV f.o. temperature
- ε 0.34 ± 0.01 eccentricity
- u_t^2/b -0.34 ± 0.07 ($b < 0$) transv. flow/temp. grad
- χ^2 256 (66 with theory error)

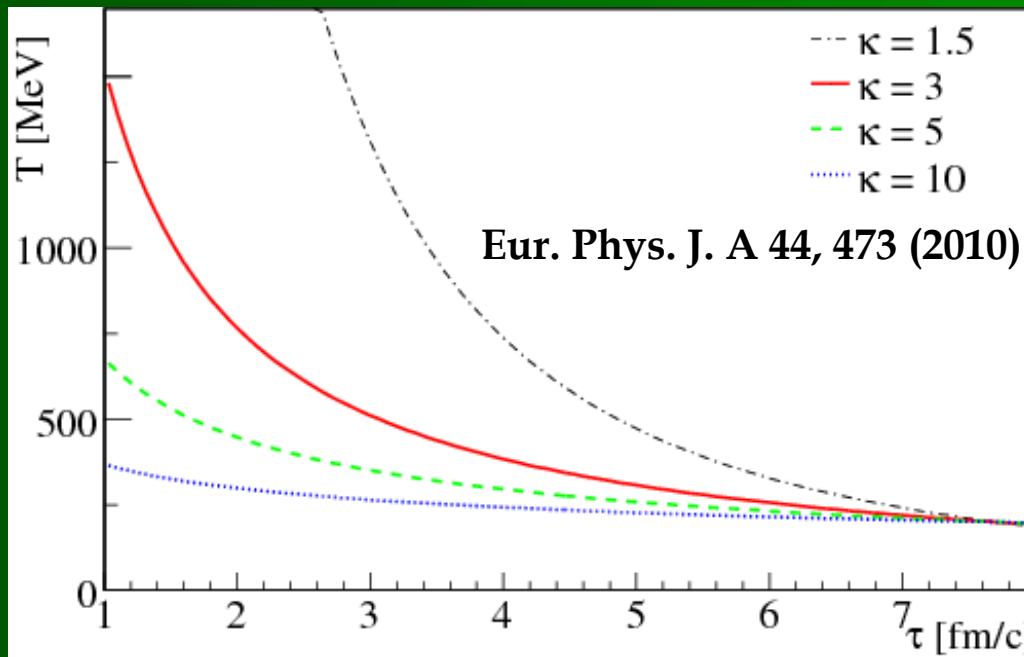
Direct photon spectrum

- Fits to 0-92% centrality PHENIX data
- Parameters from hadronic fit
- Important new parameter: $\kappa=7.7\pm0.8$



Implications on initial temperature

- From hadronic observables:

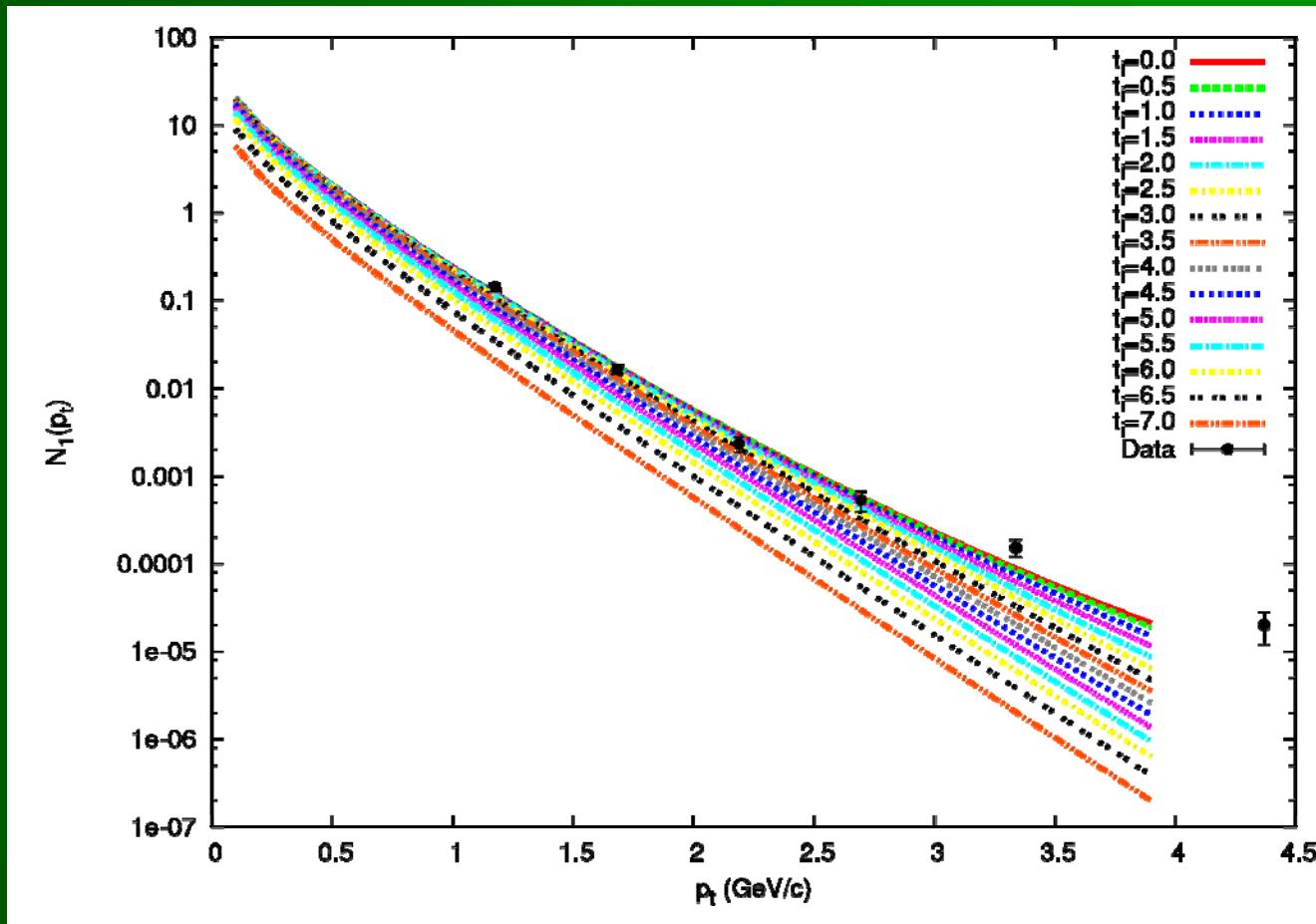


- EoS from photon spectra: $\kappa=7.7\pm0.8$
- Initial temperature (at $\tau=1$ fm/c)

$$T_i \approx 440 \text{ MeV}$$

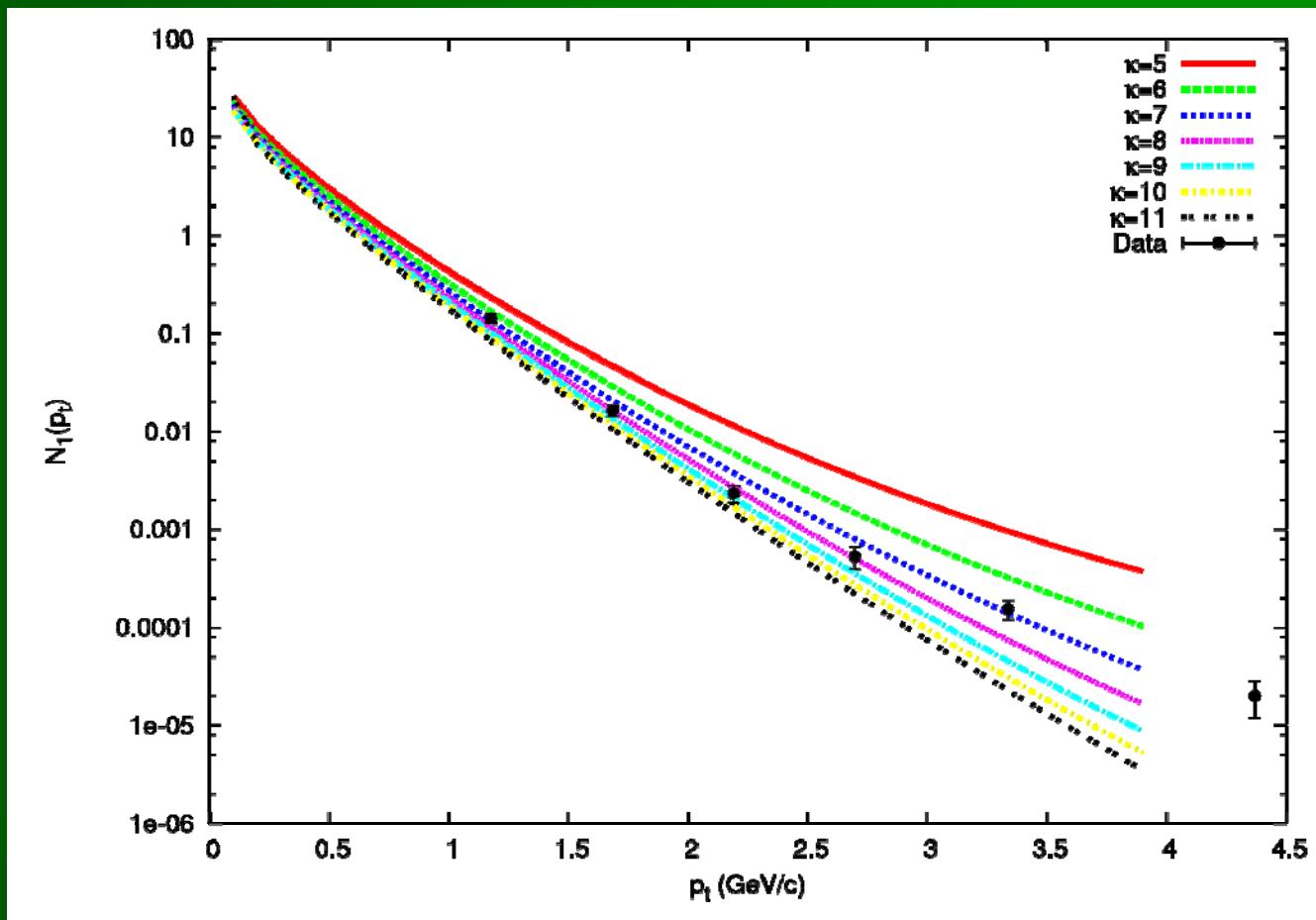
Insensitivity to the initial time

- Initial time period: small contribution

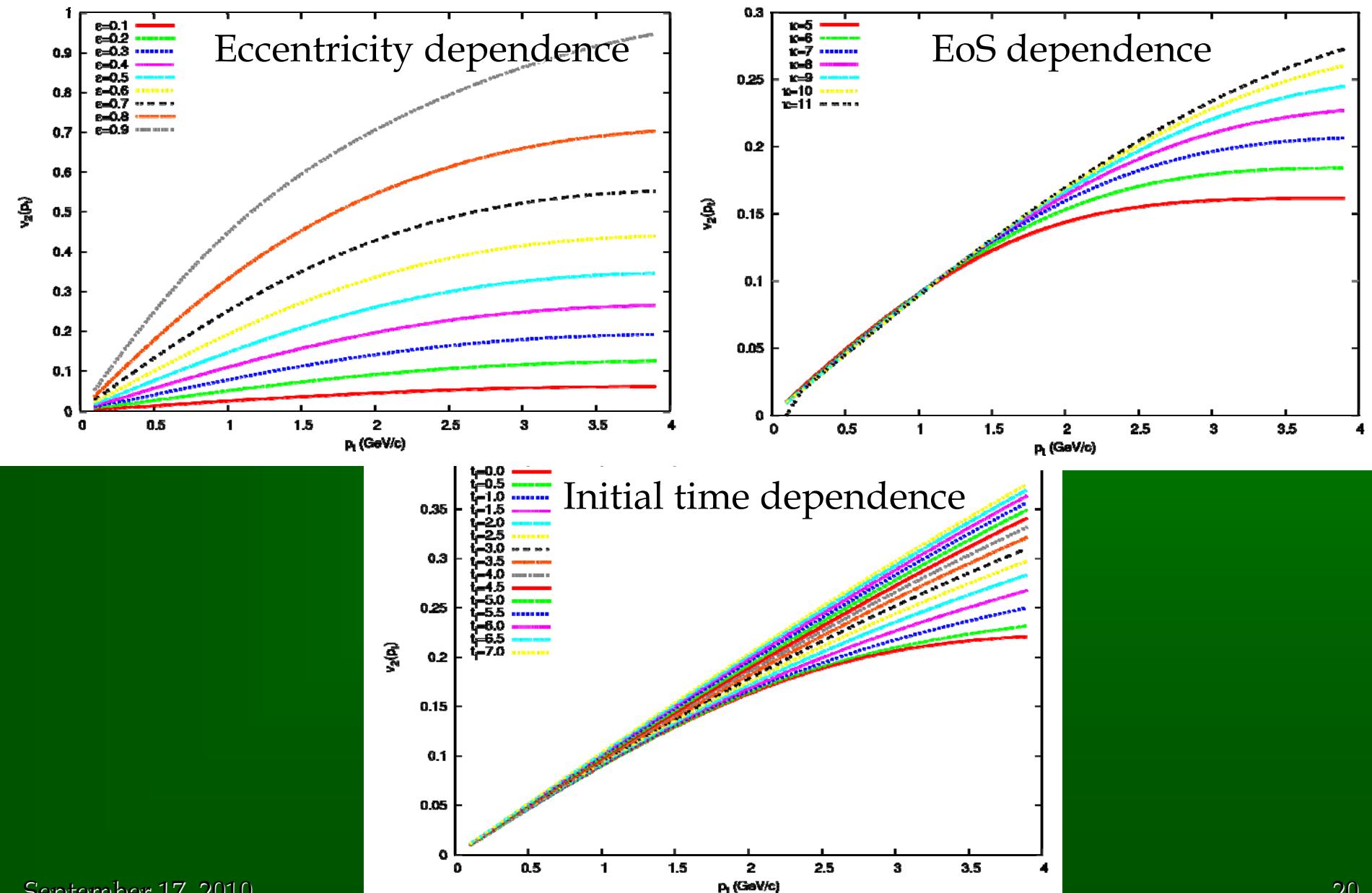


EoS dependence

- Sensitive to κ with these level of errors

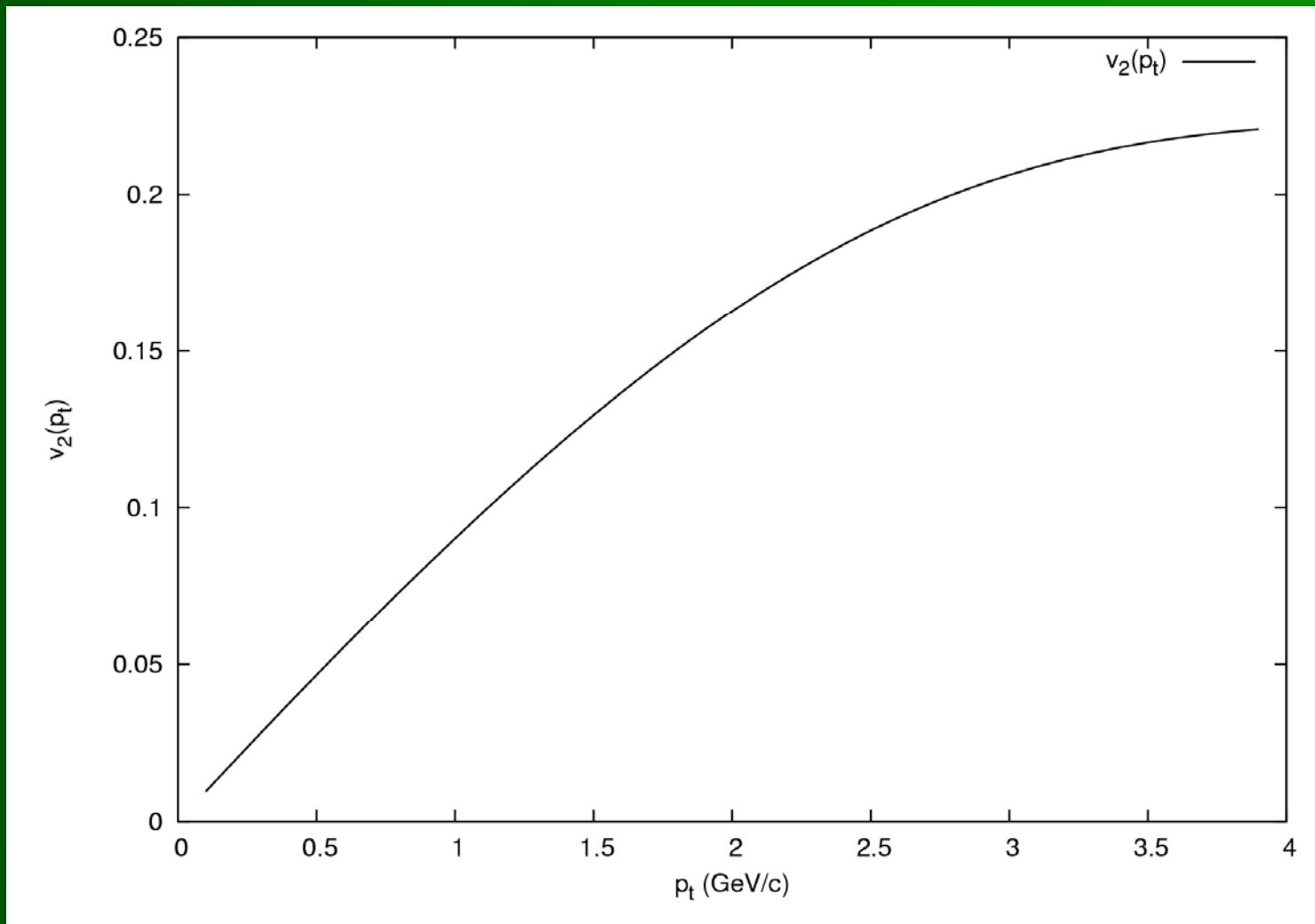


Photon elliptic flow analysis



Photon elliptic flow prediction

- With the same parameters, not fine tuned



Summary

- Revival of interest in perfect hydro
- Our model: 3+1d relativistic model w/o acceleration
- Calculated hadronic source $\rightarrow N_1, v_2, \text{HBT}$
- Calculated photon source $\rightarrow N_1, v_2$
- Photon HBT to be calculated yet
- Compared successfully to data, $\kappa=7.7\pm0.8$

$$T_i \approx 440 \text{ MeV}$$

- Predicted photon elliptic flow

**Thank you for your
attention**

Two-particle correlation radii

- Definition:

$$C_2(p_1, p_2) = \frac{N_2(p_1, p_2)}{N_1(p_1)N_1(p_2)}$$

- From the source function: $C_2(q, K) = 1 + \lambda \left| \frac{\tilde{S}(q, K)}{\tilde{S}(0, K)} \right|^2$
- Changing coordinates

$$q = p_1 - p_2, \quad K = 0.5(p_1 + p_2) \Rightarrow C_2(q, K) = 1 + \lambda \exp(-R_x^2 q_x^2 - R_y^2 q_y^2 - R_z^2 q_z^2)$$

- Result: $R_i = \frac{T_0 \tau_0 (T_i - T_0)}{m_t T_i}$

- The usual scaling (same for kaons!): $R_i^2 \sim \frac{1}{m_t}$

- Bertsch-Pratt coordinates: $R_{out} = R_{side} = 0.5(R_x^2 + R_y^2)^{1/2}$

- Freeze-out: $\tau = \text{const.} \leftrightarrow \Delta\tau = 0 \rightarrow R_{out} = R_{side}$

Summary of the fit results

- Freeze-out temperature around 200 MeV
- Flow strongly depends on centrality
- Momentum space eccentricity: 0.3-0.9
 - This is flow assymmetry
- Average transverse flow and temp. gradient:
 - Strongly coupled, ratio around 0.3-1.0 (with $b < 0$)
- Confidence levels very low
- With estimated 3% theory error: acceptable

Famous solutions

- Landau's solution (1D, developed for p+p):
 - Accelerating, implicit, complicated, 1D
 - L.D. Landau, Izv. Acad. Nauk SSSR 81 (1953) 51
 - I.M. Khalatnikov, Zhur. Eksp.Teor.Fiz. 27 (1954) 529
 - L.D. Landau and S.Z. Belenkij, Usp. Fiz. Nauk 56 (1955) 309
- Hwa-Bjorken solution:
 - Non-accelerating, explicit, simple, 1D, boost-invariant
 - R.C. Hwa, Phys. Rev. D10, 2260 (1974)
 - J.D. Bjorken, Phys. Rev. D27, 40(1983)
- Others
 - Chiu, Sudarshan and Wang
 - Baym, Friman, Blaizot, Soyeur and Czyz
 - Srivastava, Alam, Chakrabarty, Raha and Sinha

Nonrelativistic solutions

Solution	Symmetry	Density prof.	EoS	Observables
Csizmadia et al. Phys. Lett. B443:21-25, 1998	Sphere	Gaussian	$\varepsilon = \frac{3}{2} nT$	Calculated
Csörgő Central Eur.J.Phys.2: 556-565, 2004	Sphere	Arbitrary	$\varepsilon = \frac{3}{2} nT$	Not calculated
Akkelin et al. Phys.Rev. C67,2003	Ellipsoid	Gaussian ($T=T(t)$)	$\varepsilon = \kappa(T)nT$	Calculated
Csörgő Acta Phys.Polon. B37:483-494,2006	Ellipsoid	Arbitrary ($T=T(r,t)$)	$\varepsilon = \kappa nT$	Not calculated
Csörgő, Zimányi Heavy Ion Phys.17:281-293,2003	Ellipsoid	Gaussian	$\varepsilon = \kappa_e nT - B$	Calculated

Where we are

- Revival of interest, new solutions
 - Sinyukov, Karpenko, nucl-th/0505041
 - Pratt, nucl-th/0612010
 - Bialas et al.: Phys.Rev.C76:054901,2007
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 - Liao et al.: arXiv/09092284 and Phys.Rev.C80:034904,2009
 - Mizoguchi et al.: Eur.Phys.J.A40:99-108,2009
 - Beuf et al.:Phys.Rev.C78:064909,2008 (dS/dy as well!)
- Need for solutions that are:
 - accelerating + relativistic+ 3 dimensional
 - explicit + simple + compatible with the data
- Buda-Lund type of solutions: each, but not simultaneously
- Buda-Lund interpolator: hydro inspired source function, interpolates between 3-dimensional B-L solutions:

Non-relativistic, accelerating, 3d

B-L interpolator

Relativistic, non-accelerating 3d